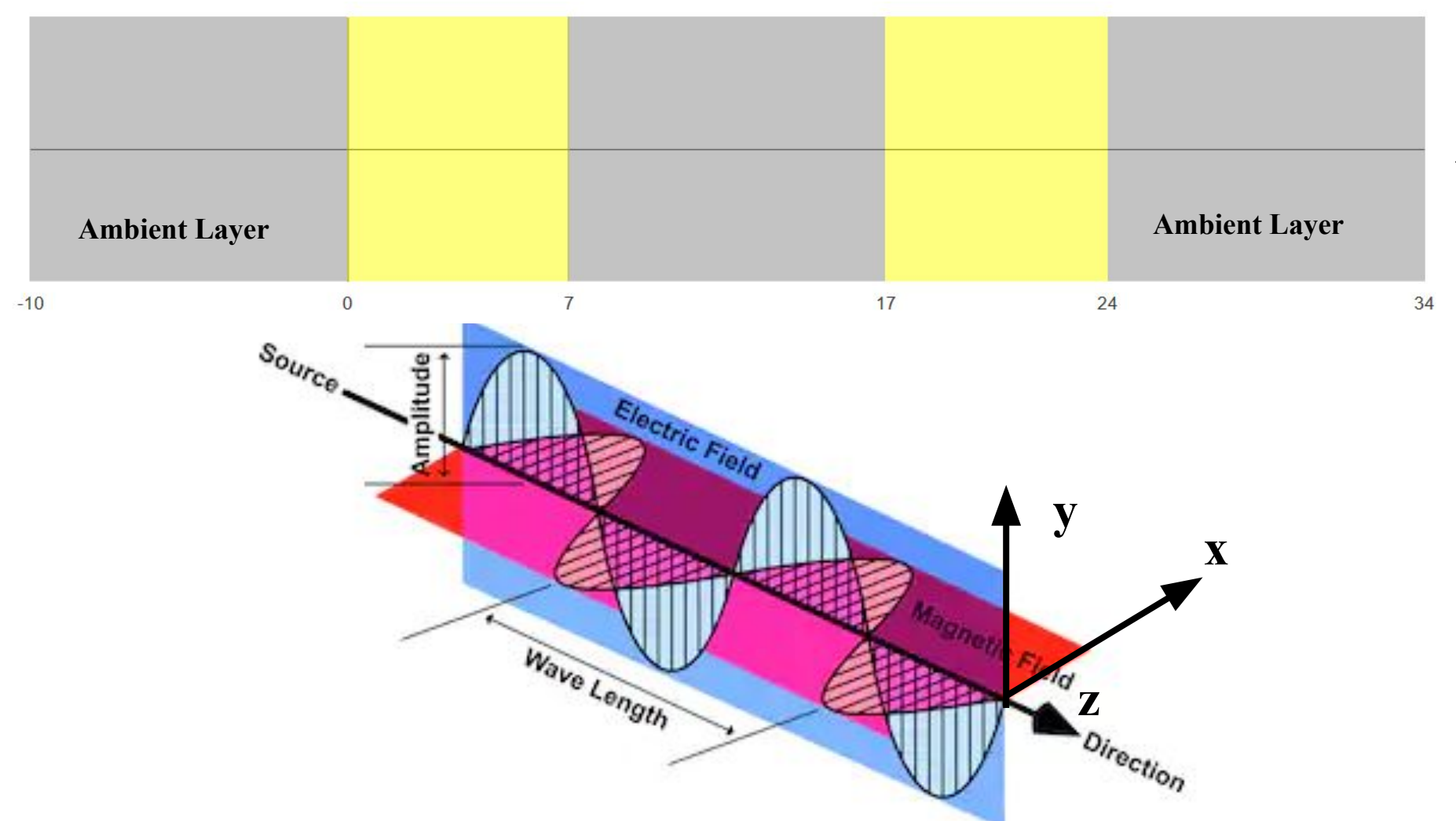


Abstract

The goal is to make available to the community a versatile, online application for the computation of electromagnetic fields in media with any number of layers having arbitrary electric and magnetic tensors. The objectives are:

- (1) to allow scientists to explore phenomena of scattering, guided modes, and resonance in the most general EM layered media
- (2) to provide a pedagogical tool for students and professionals to learn EM in layered media.



Harmonic Fields in Layered Media

Layered Structures: We can create a structure made out of any number of layers. The medium in each layer can be represented by μ and ϵ matrices.

$$\epsilon = \epsilon^{(1)} + i\epsilon^{(2)}$$

$$\epsilon = \begin{bmatrix} \epsilon_{11}^{(1)} & \epsilon_{12}^{(1)} & \epsilon_{13}^{(1)} \\ \epsilon_{12}^{(1)} & \epsilon_{22}^{(1)} & \epsilon_{23}^{(1)} \\ \epsilon_{13}^{(1)} & \epsilon_{23}^{(1)} & \epsilon_{33}^{(1)} \end{bmatrix} + i \begin{bmatrix} \epsilon_{11}^{(2)} & \epsilon_{12}^{(2)} & \epsilon_{13}^{(2)} \\ \epsilon_{12}^{(2)} & \epsilon_{22}^{(2)} & \epsilon_{23}^{(2)} \\ \epsilon_{13}^{(2)} & \epsilon_{23}^{(2)} & \epsilon_{33}^{(2)} \end{bmatrix}$$

$\epsilon^{(1)}$ and $\epsilon^{(2)}$ are Hermitian Matrices, that is that the matrix is equal to its own conjugate transpose. The same is true for μ .

$$\mu = \mu^{(1)} + i\mu^{(2)}$$

μ is a magnetic tensor and ϵ is an electric tensor.

Harmonic Field: Each layer will have an electromagnetic field, denoted by $E(x,y,z,t)$ for the electric component and $H(x,y,z,t)$ for the magnetic component. A harmonic field results when there is oscillation in the x , y , and t directions. This is controlled by the wave vector $\langle k_1, k_2 \rangle$ parallel to the layers, and the frequency ω like so

$$E = E(z)e^{i(k_1x + k_2y - \omega t)} \quad H = H(z)e^{i(k_1x + k_2y - \omega t)}$$

Scattering Problem: Using the structure built above, we can simulate physical experiments. The experiment consists of a wave source located in each of the ambient layers. The angle of incidence of these fields are related to the wave vector and the frequency is controlled by ω . The resultant total electromagnetic field will be computed and displayed.

Mathematics of Waves in Layered Media

Given an n -layered structure, values of ϵ and μ for each layer, and ω and $\langle k_1, k_2 \rangle$ values, the Maxwell partial differential equations for $E(x,y,z,t)$ and $H(x,y,z,t)$ reduce to a 4×4 system of ordinary differential equations for the tangential components of $E(z)$ and $H(z)$. In each layer, the system is

$$\frac{d}{dz} F(z) = i\kappa \begin{bmatrix} -\frac{\epsilon_{31}\kappa_1}{\epsilon_{33}} - \frac{\mu_{23}\kappa_2}{\mu_{33}} & (-\frac{\epsilon_{32} + \mu_{23}}{\epsilon_{33}})\kappa_1 & \mu_{21} + \frac{\kappa_1\kappa_2}{\epsilon_{33}} - \frac{\mu_{23}\mu_{31}}{\mu_{33}} & \mu_{23} - \frac{\kappa_1^2}{\epsilon_{33}} - \frac{\mu_{23}\mu_{32}}{\mu_{33}} \\ (-\frac{\epsilon_{31}}{\epsilon_{33}} + \frac{\mu_{13}\kappa_2}{\mu_{33}}) & -\frac{\epsilon_{32}\kappa_2}{\epsilon_{33}} - \frac{\mu_{13}\kappa_1}{\mu_{33}} & -\mu_{11} + \frac{\kappa_2^2}{\epsilon_{33}} + \frac{\mu_{13}\mu_{31}}{\mu_{33}} & -\mu_{12} - \frac{\kappa_1\kappa_2}{\epsilon_{33}} + \frac{\mu_{13}\mu_{32}}{\mu_{33}} \\ -\epsilon_{21} - \frac{\kappa_1\kappa_2}{\mu_{33}} + \frac{\epsilon_{23}\epsilon_{31}}{\epsilon_{33}} & -\epsilon_{22} + \frac{\kappa_1^2}{\mu_{33}} + \frac{\epsilon_{23}\epsilon_{32}}{\epsilon_{33}} & -\frac{\mu_{31}\kappa_1}{\mu_{33}} - \frac{\epsilon_{23}\mu_{32}}{\epsilon_{33}} & (-\frac{\mu_{32}}{\mu_{33}} + \frac{\epsilon_{23}}{\epsilon_{33}})\kappa_1 \\ \epsilon_{11} - \frac{\epsilon_{13}\epsilon_{31}}{\epsilon_{33}} - \frac{\kappa_2^2}{\mu_{33}} & \epsilon_{12} - \frac{\epsilon_{13}\epsilon_{32}}{\epsilon_{33}} + \frac{\kappa_1\kappa_2}{\mu_{33}} & (-\frac{\mu_{31}}{\mu_{33}} + \frac{\epsilon_{13}}{\epsilon_{33}})\kappa_2 & -\frac{\mu_{32}\kappa_2}{\mu_{33}} - \frac{\epsilon_{13}\kappa_1}{\epsilon_{33}} \end{bmatrix} F(z)$$

where

$$\kappa = \frac{\omega}{c}, \quad \kappa_1 = \frac{k_1}{\kappa}, \quad \kappa_2 = \frac{k_2}{\kappa}, \quad c = 1, \quad \text{and} \quad F(z) = \begin{bmatrix} E_1(z) \\ E_2(z) \\ H_1(z) \\ H_2(z) \end{bmatrix}$$

The solution in each layer is a combination of four mod

$\bar{v}e^{\lambda z}$, where \bar{v} is an eigenvector and λ is an eigenvalue:

$$\text{If } F(z) = c_1\bar{v}_1e^{\lambda_1 z} + c_2\bar{v}_2e^{\lambda_2 z} + c_3\bar{v}_3e^{\lambda_3 z} + c_4\bar{v}_4e^{\lambda_4 z}$$

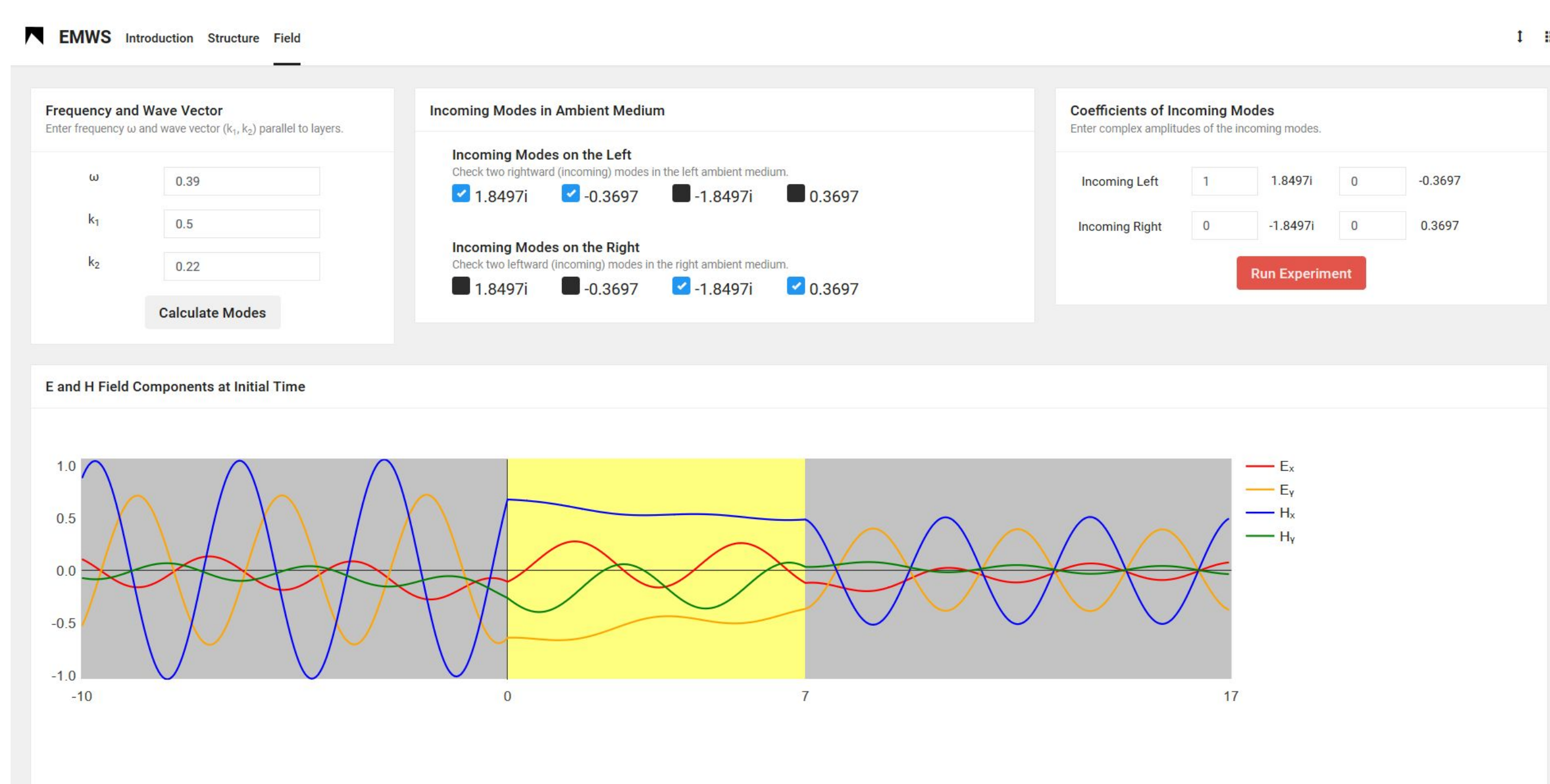
Continuity at the l^{th} interface is expressed as $n^+c^l = n^-c^{l+1}$.

This gives an $4(n-1) \times 4n$ homogeneous linear system:

$$(n=6) \quad \begin{bmatrix} n_1^+ & -n_2^- & 0 & 0 & 0 & 0 \\ 0 & n_2^- & -n_3^- & 0 & 0 & 0 \\ 0 & 0 & n_3^- & -n_4^- & 0 & 0 \\ 0 & 0 & 0 & n_4^- & -n_5^- & 0 \\ 0 & 0 & 0 & 0 & n_5^- & -n_6^- \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \\ c_3 \\ c_4 \\ c_5 \\ c_6 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \end{bmatrix}$$

For the n^{th} layer, this results in a nonhomogeneous $4n \times 4n$ linear system. By solving for all the coefficients we obtain the full field in the structure.

Graphical User Interface

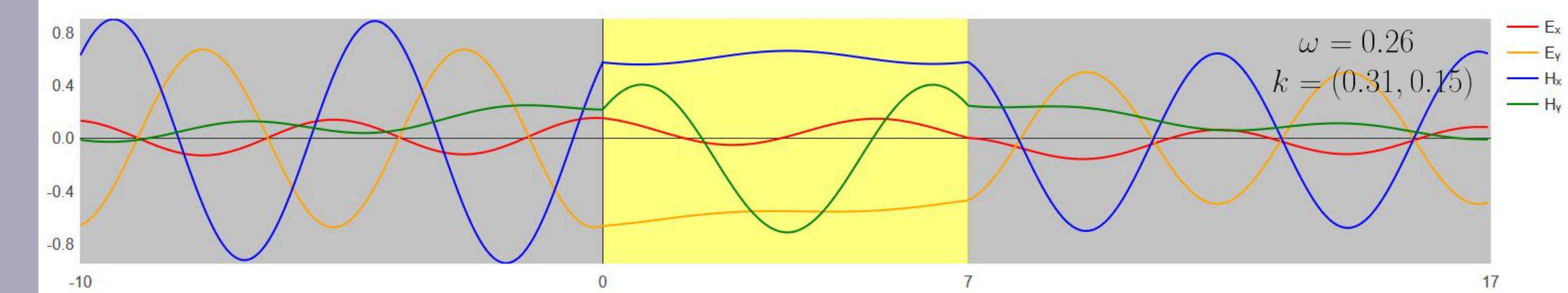
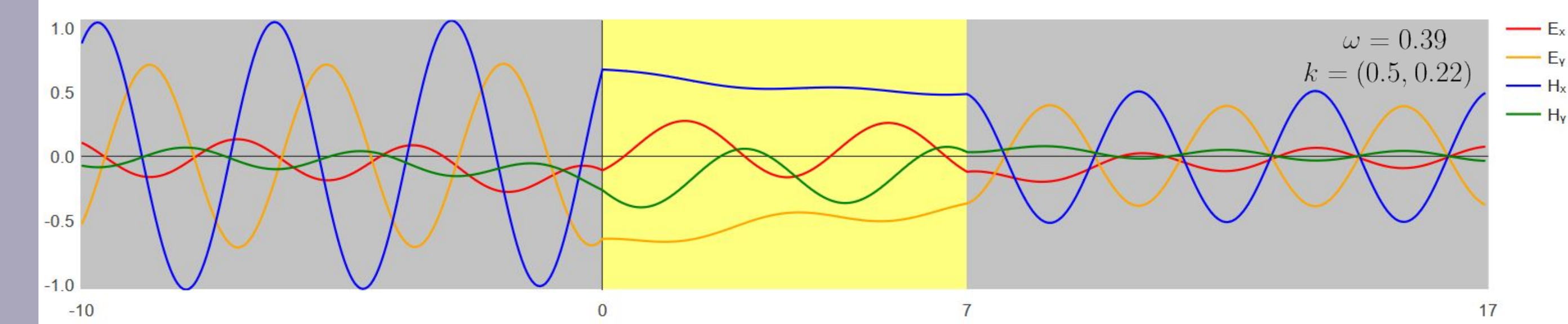


Screenshot of current graphical user interface (GUI)

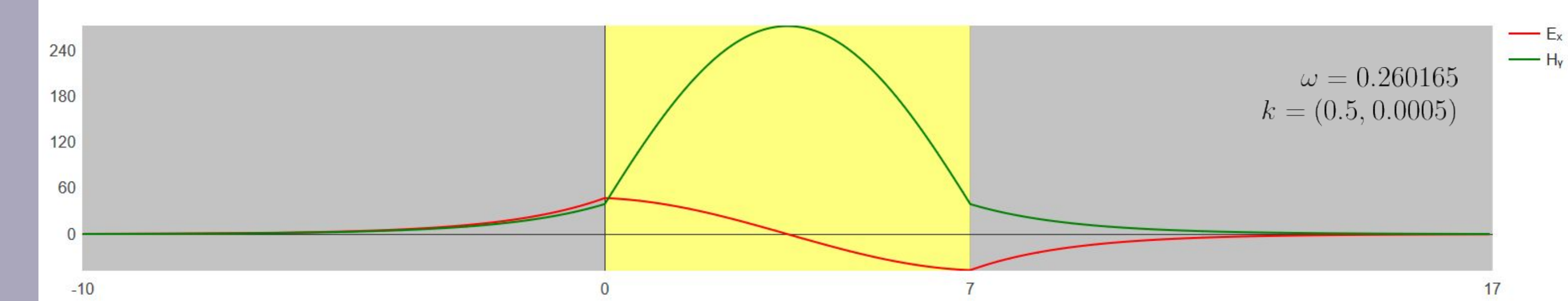
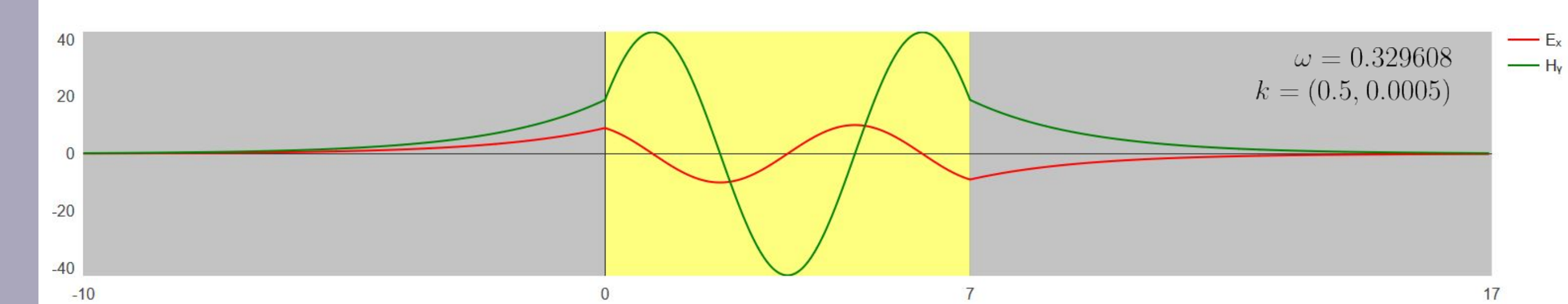
<https://www.math.lsu.edu/~shipman/EMWS/html/dashboard.4.html>

Simulations

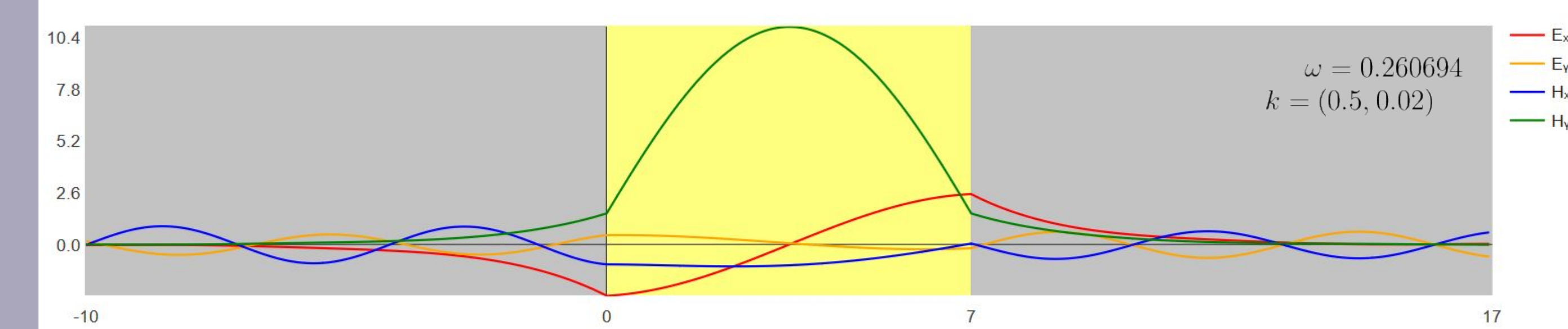
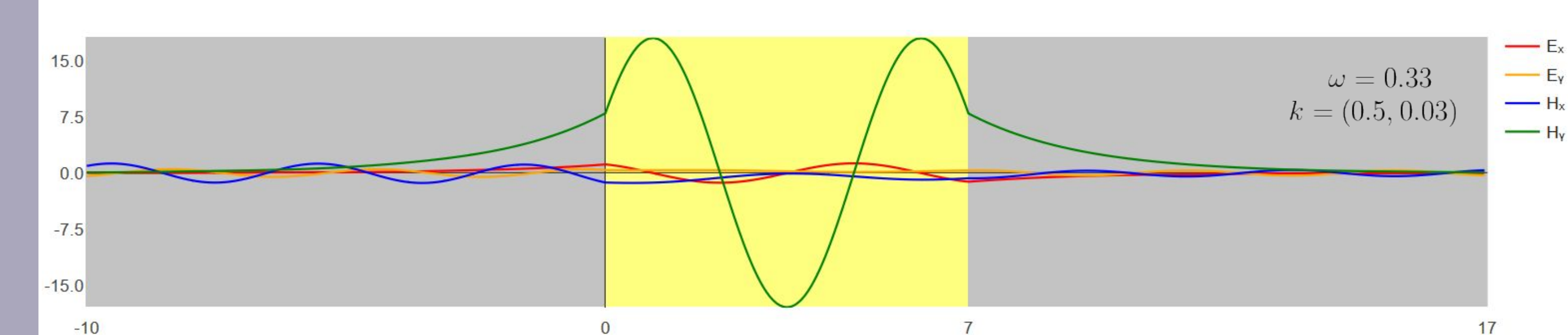
• Scattering



• Guided Modes



• Resonance



Acknowledgements

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