

The variational approach to homogenization

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I. The stochastic homogenization error

The coefficient field: the map $a: \mathbb{R}^d \times \Omega \rightarrow \mathbb{R}^{d \times d}$ satisfies

$$\xi \cdot a \xi \geq \lambda |\xi|^2 \quad \text{and} \quad \xi \cdot a^{-1} \xi \geq \Lambda^{-1} |\xi|^2,$$

and, for a measure-preserving transformation group $\{\tau_x: \Omega \rightarrow \Omega\}_{x \in \mathbb{R}^d}$,

$$a(x + y, \omega) = a(y, \tau_x \omega) \quad \text{for all } x, y \in \mathbb{R}^d \text{ and } \omega \in \Omega.$$

Furthermore, the transformation group is ergodic:

$$f(\omega) = f(\tau_x \omega) \in L^\infty(\Omega) \quad \text{for all } x \in \mathbb{R}^d \text{ if and only if } f \text{ is constant.}$$

The equation: for a fixed environment $\omega \in \Omega$, and $U \subseteq \mathbb{R}^d$,

$$-\nabla \cdot a^\varepsilon \nabla u_\varepsilon = f \quad \text{in } U \quad \text{with } u_\varepsilon = 0 \quad \text{on } \partial U.$$

Qualitative homogenization: there exists $\bar{a} \in \mathbb{R}^{d \times d}$ such that, $\mathbb{P}$ -a.s.,

$$u_\varepsilon \rightharpoonup \bar{u} \quad \text{weakly in } H^1(U) \quad \text{and} \quad u_\varepsilon \rightarrow \bar{u} \quad \text{strongly in } L^2(U),$$

for $\bar{u}$ solving

$$-\nabla \cdot \bar{a} \nabla \bar{u} = f \quad \text{in } U \quad \text{with } \bar{u} = 0 \quad \text{on } \partial U.$$

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The homogenization corrector: there $\mathbb{P}$ -a.s. exist $\phi_i \in H^1_{\text{loc}}(\mathbb{R}^d)$ satisfying

$$-\nabla \cdot a(\nabla \phi_i + e_i) = 0 \text{ in } \mathbb{R}^d \text{ with } \nabla \phi_i \text{ stationary.}$$

That is, there exists $D\Phi_i \in L^2_{\text{pot}}(\Omega)$ such that

$$\nabla \phi_i(x, \omega) = D\Phi_i(\tau_x \omega).$$

Sublinearity: for $\mathbb{P}$ -a.e. $\omega \in \Omega$, for every $R \in (0, \infty)$,

$$\lim_{\varepsilon \rightarrow 0} \|\varepsilon \phi_i^\varepsilon\|_{L^2(B_R)} = \lim_{r \rightarrow \infty} \frac{1}{r} \left(\int_{B_{rR}} |\phi_i|^2 \right) = 0.$$

The homogenized coefficient: as in the periodic case,

$$\bar{a}e_i = \langle a(e_i + \nabla \phi_i) \rangle,$$

is uniformly elliptic (with the same constants) and symmetric (if a is).

The two-scale expansion: we define

$$u_\varepsilon \simeq \tilde{u}_\varepsilon = \bar{u} + \varepsilon \phi_i^\varepsilon \partial_i \bar{u},$$

for $\bar{u}$ solving

$$-\nabla \cdot \bar{a} \nabla \bar{u} = f \text{ in } U \text{ with } \bar{u} = 0 \text{ on } \partial U.$$

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The two-scale expansion: we define

$$u_\varepsilon \simeq \tilde{u}_\varepsilon = \bar{u} + \varepsilon \phi_i^\varepsilon \partial_i \bar{u}.$$

The stochastic homogenization error: we have

$$w_\varepsilon = u_\varepsilon - \bar{u} - \varepsilon \phi_i^\varepsilon \partial_i \bar{u}.$$

The equation: as in the periodic case,

$$\begin{aligned} -\nabla \cdot a^\varepsilon \nabla w_\varepsilon &= (a^\varepsilon (e_i + \nabla \phi_i) - \bar{a} e_i) \cdot \nabla \partial_i \bar{u} + \nabla \cdot (\varepsilon \phi_i^\varepsilon a^\varepsilon \nabla \partial_i \bar{u}) \\ &= q_i^\varepsilon \cdot \nabla \partial_i \bar{u} + \nabla \cdot (\varepsilon \phi_i^\varepsilon a^\varepsilon \nabla \partial_i \bar{u}), \end{aligned}$$

for the flux $q_i^\varepsilon = a^\varepsilon (e_i + \nabla \phi_i) - \bar{a} e_i$.

The flux correction: in the periodic case,

$$-\Delta \sigma_{ijk} = \partial_j q_{ik} - \partial_k q_{ij} \quad \text{and} \quad \langle \sigma_{ijk} \rangle_{\mathbb{T}^d} = 0,$$

which, when lifted to Ω , for $D_i \Psi = \lim_{h \rightarrow 0} h^{-1} (\Psi(\tau_{he_i} \omega) - \Psi(\omega))$,

$$-D \cdot D \Sigma_{ijk} = D_j Q_{ik} - D_k Q_{ij} \quad \text{in } L^2_{\text{pot}}(\Omega),$$

for $q_i(x, \omega) = Q_i(\tau_x \omega)$, and in the weak form

$$\mathbb{E}[D\Psi \cdot D\Sigma_{ijk}] = \mathbb{E}[Q_{ik} D_j \Psi - Q_{ij} D_k \Psi] \quad \text{for all } D\Psi \in L^2_{\text{pot}}(\Omega).$$

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The flux correction: there exists a unique $D\Sigma_{ijk} \in L^2_{\text{pot}}(\Omega)$ satisfying

$$\mathbb{E}[D\Psi \cdot D\Sigma_{ijk}] = \mathbb{E}[Q_{ik}D_j\Psi - Q_{ij}D_k\Psi] \quad \text{for all } D\Psi \in L^2_{\text{pot}}(\Omega).$$

As with the correctors, there $\mathbb{P}$ -a.s. exist $\sigma_{ijk} \in H^1_{\text{loc}}(\mathbb{R}^d)$ satisfying

$$-\Delta\sigma_{ijk} = \partial_j q_{ik} - \partial_k q_{ij} \quad \text{in } \mathbb{R}^d \text{ with } \nabla\sigma_{ijk} \text{ stationary.}$$

Sublinearity: since the gradient is stationary with $\mathbb{E}[\nabla\sigma_{ijk}] = 0$,

$$\lim_{\varepsilon \rightarrow 0} \|\varepsilon\sigma_{ijk}^\varepsilon\|_{L^2(B_R)} = \lim_{r \rightarrow \infty} \frac{1}{r} \left(\int_{B_{rR}} |\sigma_{ijk}|^2 \right) = 0.$$

The equation: $\mathbb{P}$ -a.s. for the skew-symmetric matrix $\sigma_i = (\sigma_{ijk})$, using $\nabla \cdot q_i = 0$,

$$-\Delta((\nabla \cdot \sigma_i)_j - q_{ij}) = 0 \quad \text{in } \mathbb{R}^d,$$

which implies $(\nabla \cdot \sigma_i)_j - q_{ij}$ is constant, and by the ergodic theorem

$$\lim_{r \rightarrow \infty} \int_{B_r} ((\nabla \cdot \sigma_i)_j - q_{ij}) = \mathbb{E}[(\nabla \cdot \sigma_i)_j - q_{ij}] = 0.$$

Therefore,

$$\nabla \cdot \sigma_i = q_i \quad \text{for } (\nabla \cdot \sigma_i)_j = \partial_k \sigma_{ijk}.$$

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The stochastic homogenization error: we have

$$w_\varepsilon = u_\varepsilon - \bar{u} - \varepsilon \phi_i^\varepsilon \partial_i \bar{u}.$$

The equation: as in the periodic case, using the skew-symmetry of σ_i ,

$$\begin{aligned} -\nabla \cdot a^\varepsilon \nabla w_\varepsilon &= q_i^\varepsilon \cdot \nabla \partial_i \bar{u} + \nabla \cdot (\varepsilon \phi_i^\varepsilon a^\varepsilon \nabla \partial_i \bar{u}) \\ &= (\nabla \cdot \sigma_i^\varepsilon) \cdot \nabla \partial_i \bar{u} + \nabla \cdot (\varepsilon \phi_i^\varepsilon a^\varepsilon \nabla \partial_i \bar{u}) \\ &= \nabla \cdot ((\varepsilon \phi_i^\varepsilon a^\varepsilon - \varepsilon \sigma_i^\varepsilon) \nabla \partial_i \bar{u}). \end{aligned}$$

The boundary layer: for $0 < r \ll 1$ let

$$U_r = \{x \in U : d(x, \partial U) \geq r\}$$

and for $\eta_r = 1$ on U_r , $\eta_r = 0$ on $\mathbb{R}^d \setminus U_{r/2}^\circ$, and $|\nabla \eta_r| \leq c/r$,

$$w_{\varepsilon,r} = u_\varepsilon - \bar{u} - \varepsilon \phi_i^\varepsilon \eta_r \partial_i \bar{u} \in H_0^1(U).$$

The equation: we have that

$$-\nabla \cdot a^\varepsilon \nabla w_{\varepsilon,r} = \nabla \cdot ((1 - \eta_r)(a^\varepsilon - \bar{a}) \nabla \bar{u}) + \nabla \cdot ((\varepsilon \phi_i^\varepsilon a^\varepsilon - \varepsilon \sigma_i^\varepsilon) \nabla (\eta_r \partial_i \bar{u})).$$

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The boundary expansion: since $w_{\varepsilon,r} = u_\varepsilon - \bar{u} - \varepsilon \eta_r \phi_i^\varepsilon \partial_i \bar{u} \in H_0^1(U)$,

$$-\nabla \cdot a^\varepsilon \nabla w_{\varepsilon,r} = \nabla \cdot ((1 - \eta_r)(a^\varepsilon - \bar{a}) \nabla \bar{u}) + \nabla \cdot ((\varepsilon \phi_i^\varepsilon a^\varepsilon - \varepsilon \sigma_i^\varepsilon) \nabla (\eta_r \partial_i \bar{u})).$$

The energy estimate: for $c \in (0, \infty)$ depending on λ , Λ , and U ,

$$\begin{aligned} \int_U |\nabla w_{\varepsilon,r}|^2 &\leq cr \|\nabla^2 \bar{u}\|_{L^\infty(U \setminus U_r)}^2 + cr^{-2} \|\nabla \bar{u}\|_{L^\infty(U_r \setminus U_{r/2})}^2 \|\varepsilon(\phi_i^\varepsilon, \sigma_i^\varepsilon)\|_{L^2(U_r \setminus U_{r/2})}^2 \\ &\quad + c \|\nabla^2 \bar{u}\|_{L^\infty(U_{r/2})}^2 \|\varepsilon(\phi_i^\varepsilon, \sigma_i^\varepsilon)\|_{L^2(U_{r/2})}^2. \end{aligned}$$

The full two-scale expansion: for $w_\varepsilon = u_\varepsilon - \bar{u} - \varepsilon \phi_i^\varepsilon \partial_i \bar{u}$, for $r = \varepsilon$,

$$\begin{aligned} \|\nabla w_\varepsilon - \nabla w_{\varepsilon,r}\|_{L^2(U)}^2 &= \|\nabla(\varepsilon(1 - \eta_r) \phi_i^\varepsilon \partial_i \bar{u})\|_{L^2(U)}^2 \\ &\leq c(\varepsilon^{-1} \|\nabla \bar{u}\|_{L^\infty(U \setminus U_r)}^2 \|\varepsilon \phi_i^\varepsilon\|_{L^2(U \setminus U_r)}^2 + \|\nabla \bar{u}\|_{L^\infty(U \setminus U_r)}^2 \|\nabla \phi_i^\varepsilon\|_{L^2(U \setminus U_r)}^2) \\ &\quad + c\varepsilon \|\nabla^2 \bar{u}\|_{L^\infty(U \setminus U_r)}^2 \|\phi_i^\varepsilon\|_{L^2(U \setminus U_r)}^2. \end{aligned}$$

Qualitative homogenization: if $\bar{u} \in W^{2,\infty}(U)$,

$$\limsup_{\varepsilon \rightarrow 0} \int_U |\nabla w_\varepsilon|^2 = 0.$$

I. The stochastic homogenization error

Qualitative stochastic homogenization

Let $U \subseteq \mathbb{R}^d$ be a smooth, bounded domain, let $f \in C^\alpha(\overline{U})$ for some $\alpha \in (0, 1)$, and let $a: \mathbb{R}^d \times \Omega \rightarrow \mathbb{R}^{d \times d}$ be a uniformly elliptic, stationary, and ergodic diffusion matrix. There $\mathbb{P}$ -a.s. exist homogenization correctors $\phi_i \in H_{\text{loc}}^1(\mathbb{R}^d)$ satisfying

$$-\nabla \cdot a(e_i + \nabla \phi_i) = 0 \text{ in } \mathbb{R}^d \text{ with } \nabla \phi_i \text{ stationary,}$$

that define the homogenized matrix

$$\bar{a}e_i = \langle a(e_i + \nabla \phi_i) \rangle,$$

and the macroscopic limit $\bar{u}$ solving

$$-\nabla \cdot \bar{a} \nabla \bar{u} = f \text{ in } U \text{ with } \bar{u} = 0 \text{ on } \partial U.$$

The stochastic homogenization error

$$w_\varepsilon = u_\varepsilon - \bar{u} - \varepsilon \phi_i^\varepsilon \partial_i \bar{u},$$

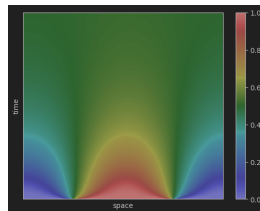
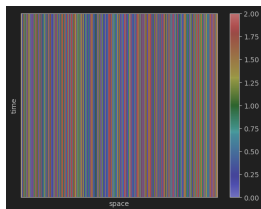
satisfies

$$\lim_{\varepsilon \rightarrow 0} \int_U |\nabla w_\varepsilon|^2 = 0 \text{ and, so, } \nabla u_\varepsilon \simeq \nabla \bar{u} + \partial_i \bar{u} \nabla \phi_i^\varepsilon.$$

I. The stochastic homogenization error

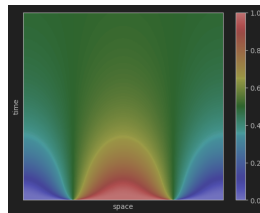
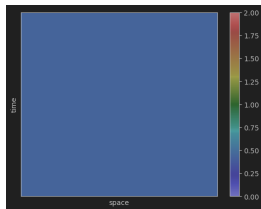
The equation: for $\varepsilon = 1/1000$ in $\mathbb{T}^1 \times (0, \infty)$,

$$\partial_t u_\varepsilon = (1/2) \nabla \cdot a^\varepsilon \nabla u_\varepsilon \text{ in } \mathbb{T}^1 \times (0, \infty) \text{ with } u_0 = \mathbf{1}_{\{1/4, 3/4\}}.$$



A macroscopic limit: choosing $\bar{a} = \langle a^{-1} \rangle^{-1} \simeq 0.43588$,

$$\partial_t \bar{u} = (1/2) \nabla \cdot \bar{a} \nabla \bar{u} \text{ in } \mathbb{T}^1 \times (0, \infty) \text{ with } u_0 = \mathbf{1}_{\{1/4, 3/4\}}.$$



II. The variational approach

The equation: for a symmetric diffusion matrix a , the solution

$$-\nabla \cdot a^\varepsilon \nabla u_\varepsilon = f \text{ in } U \text{ with } u_\varepsilon = 0 \text{ on } \partial U,$$

satisfies

$$u_\varepsilon = \operatorname{argmin}_{v \in H_0^1(U)} \int_U (\nabla v \cdot a^\varepsilon \nabla v - v f).$$

The approximate gradient: by the two-scale expansion,

$$\nabla u_\varepsilon \simeq \nabla \bar{u} + \partial_i \bar{u} \nabla \phi_i^\varepsilon = (e_i + \nabla \phi_i^\varepsilon) \partial_i \bar{u},$$

and, therefore,

$$\int_U \nabla v \cdot a^\varepsilon \nabla v \simeq \sum_{ij} \int_U ((e_i + \nabla \phi_i^\varepsilon) \cdot a^\varepsilon (e_j + \nabla \phi_j^\varepsilon)) \partial_i \bar{u} \partial_j \bar{u}.$$

The effective energy: we have that

$$\frac{1}{|U|} \int_U (e_i + \nabla \phi_i^\varepsilon) \cdot a^\varepsilon (e_j + \nabla \phi_j^\varepsilon) = \oint_{U/\varepsilon} (e_i + \nabla \phi_i) \cdot a (e_j + \nabla \phi_j).$$

Polarization: we will analyze on large-scales, for $\xi \in \mathbb{R}^d$,

$$\mu(U, \xi) = \frac{1}{2} \int_U (\xi + \nabla \phi_\xi) \cdot a (\xi + \nabla \phi_\xi),$$

for $\phi_\xi = \xi_i \phi_i$ satisfying $-\nabla \cdot a (\xi + \nabla \phi_\xi) = 0$.

II. The variational approach

The energy: for every $V \subseteq \mathbb{R}^d$, for $\xi \in \mathbb{R}^d$, and for $\ell_\xi(x) = x \cdot \xi$,

$$\mu(V, \xi) = \inf_{v \in \ell_\xi + H_0^1(U)} \frac{1}{2} \int_V \nabla v \cdot a \nabla v.$$

The Euler–Lagrange equation: the minimizer $v_{\xi,V}$ satisfies

$$-\nabla \cdot a \nabla v_{\xi,V} = 0 \text{ in } V \text{ with } v_{\xi,V} = \ell_\xi \text{ on } \partial V.$$

The finite-volume corrector: if $\phi_{\xi,V} = v_{\xi,V} - \ell_\xi$,

$$-\nabla \cdot a(\nabla \phi_{\xi,V} + \xi) = 0 \text{ in } V \text{ with } \phi_{\xi,V} = 0 \text{ on } \partial V.$$

Approximate correctors: for $\alpha \in (0, 1)$,

$$\alpha \phi_{i,\alpha} - \nabla \cdot a(e_i + \nabla \phi_{i,\alpha}) = 0 \text{ in } \mathbb{R}^d,$$

satisfy that

$$\phi_{i,\alpha}(x) = \int_0^\infty e^{-\alpha t} \rho_i(x, t) dt,$$

for

$$\partial_t \rho_i = \nabla \cdot a \nabla \rho_i \text{ in } \mathbb{R}^d \times (0, \infty) \text{ with } \rho_i(\cdot, 0) = \nabla \cdot (a e_i).$$

II. The variational approach

The energy: for every $V \subseteq \mathbb{R}^d$, for $\xi \in \mathbb{R}^d$, and for $\ell_\xi(x) = x \cdot \xi$,

$$\mu(V, \xi) = \inf_{v \in \ell_\xi + H_0^1(V)} \frac{1}{2} \int_V \nabla v \cdot a \nabla v = \frac{1}{2} \int_V (\xi + \nabla \phi_{\xi, V}) \cdot a (\xi + \nabla \phi_{\xi, V}),$$

for $-\nabla \cdot a(\xi + \nabla \phi_{\xi, V}) = 0$ in V with $\phi_{\xi, V} = 0$ on ∂V .

The coarse-grained matrix: since the map $\xi \rightarrow (\xi + \nabla \phi_{\xi, V})$ is linear,

the map $\xi \rightarrow \mu(V, \xi)$ is quadratic in $\xi \in \mathbb{R}^d$.

There exists a symmetric $a(V) \in \mathbb{R}^{d \times d}$ such that

$$\mu(V, \xi) = \frac{1}{2} (\xi \cdot a(V) \xi) \quad \text{for every } \xi \in \mathbb{R}^d.$$

Subadditivity: if $V = \coprod_{i=1}^N V_i$ then

$$|V| \mu(V, \xi) \leq \sum_{i=1}^N |V_i| \mu(V_i, \xi) \quad \text{and, so,} \quad \mu(V, \xi) \leq \sum_{i=1}^N \frac{|V_i|}{|V|} \mu(V_i, \xi).$$

Upper bound: for every $V \subseteq \mathbb{R}^d$ and $\xi \in \mathbb{R}^d$, choosing $v = \ell_\xi$,

$$\mu(V, \xi) \leq \frac{1}{2} \left(\xi \cdot \left(\int_V a \right) \xi \right) \leq \Lambda |\xi|^2.$$

II. The variational approach

A subadditive quantity: for the set $\mathcal{V}$ of bounded, open subsets of $\mathbb{R}^d$,

$$s: \Omega \times \mathcal{V} \rightarrow \mathbb{R},$$

is subadditive if $\mathbb{P}$ -a.s. whenever $V = \coprod_{i=1}^n V_i$, we have $s(\omega, V) \leq \sum_{i=1}^n s(\omega, V_i)$.

Kingman subadditive ergodic theorem [Akcoglu, Krengel; 1981]

Let $(\Omega, \mathcal{F}, \mathbb{P})$ be a probability space with an ergodic, $\mathbb{Z}^d$ -stationary probability measure $\mathbb{P}$: for an ergodic measure-preserving transformation group $\{\tau_z : \Omega \rightarrow \Omega\}_{z \in \mathbb{Z}^d}$, for every $f \in L^\infty(\Omega)$,

$$f(\omega) = f(\tau_z \omega) \text{ in } L^\infty(\Omega) \text{ for every } z \in \mathbb{Z}^d.$$

If $s: \Omega \times \mathcal{V} \rightarrow \mathbb{R}$ is a subadditive quantity that is $\mathbb{Z}^d$ stationary:

$$s(\omega, V + z) = s(\tau_z \omega, V) \text{ for every } V \in \mathcal{V} \text{ and } z \in \mathbb{Z}^d,$$

and bounded: for some $\Lambda \in (0, \infty)$, we have $s(V) \leq \Lambda$ $\mathbb{P}$ -a.s. for every $V \in \mathcal{V}$.

Then, there exists $\bar{u} \in \mathbb{R}$ such that, $\mathbb{P}$ -a.s. for every $V \in \mathcal{V}$,

$$\lim_{t \rightarrow \infty} s(\omega, tV) = \bar{u}.$$

II. The variational approach

The energy: for every $V \subseteq \mathbb{R}^d$, for $\xi \in \mathbb{R}^d$, and for $\ell_\xi(x) = x \cdot \xi$,

$$\mu(V, \xi) = \inf_{v \in \ell_\xi + H_0^1(V)} \frac{1}{2} \int_V \nabla v \cdot a \nabla v = \frac{1}{2} \int_V (\xi + \nabla \phi_{\xi, V}) \cdot a (\xi + \nabla \phi_{\xi, V}).$$

The coarse-grained matrix: there exists a symmetric $a(V) \in \mathbb{R}^{d \times d}$ such that

$$\mu(V, \xi) = \frac{1}{2} (\xi \cdot a(V) \xi) \quad \text{for every } \xi \in \mathbb{R}^d.$$

Subadditivity: if $V = \coprod_{i=1}^N V_i$ then

$$|V| \mu(V, \xi) \leq \sum_{i=1}^N |V_i| \mu(V_i, \xi) \quad \text{and, so,} \quad \mu(V, \xi) \leq \sum_{i=1}^N \frac{|V_i|}{|V|} \mu(V_i, \xi).$$

Upper and lower bounds: for every $V \in \mathcal{V}$ and $\xi \in \mathbb{R}^d$, choosing $v = \ell_\xi$,

$$\frac{\lambda}{2} |\xi|^2 \leq \frac{\lambda}{2} \int_V |\xi + \phi_{\xi, V}|^2 \leq \mu(V, \xi) \leq \frac{1}{2} (\xi \cdot (\int_V a) \xi) \leq \frac{\Lambda}{2} |\xi|^2.$$

The subadditive ergodic theorem: for every $\xi \in \mathbb{R}^d$ and $V \in \mathcal{V}$,

$$\lim_{t \rightarrow \infty} \mu(tV, \xi) = \bar{\mu}_\xi.$$

The homogenized matrix: for a symmetric, uniformly elliptic $\bar{a} \in \mathbb{R}^{d \times d}$,

$$\lim_{t \rightarrow \infty} a(tV) = \bar{a}.$$

II. The variational approach

A useful decomposition: for every $n \in \mathbb{N}$, up to a set of Lebesgue measure zero,

$$\square_n = \left[-\frac{3^n}{2}, \frac{3^n}{2} \right] \text{ so } \mathbb{R}^d = \coprod_{z \in 3^n \mathbb{Z}^d} (z + \square_n).$$

The finite-volume correctors: for every $n \in \mathbb{N}$ and $\xi \in \mathbb{R}^d$, for the minimizer $v_{n,\xi}$,

$$-\nabla \cdot a \nabla v_{n,\xi} = 0 \text{ in } \square_n \text{ with } v_{n,\xi} = \ell_\xi \text{ on } \partial \square_n.$$

we define the finite volume corrector

$$\phi_{n,\xi} = v_{n,\xi} - \ell_\xi \text{ satisfies } -\nabla \cdot a(\xi + \nabla \phi_{n,\xi}) = 0 \text{ in } \square_n \text{ with } \phi_{n,\xi} = 0 \text{ on } \partial \square_n.$$

The finite volume flux: since we have that

$$\int_{\square_n} a(\nabla \phi_{n,\xi} + \xi) = \int_{\square_n} a \nabla v_{n,\xi} = a(\square_n) \xi.$$

the finite-volume flux is

$$q_{n,\xi} = a(\nabla \phi_{n,\xi} + \xi) - a(\square_n) \xi,$$

which satisfies, $\mathbb{P}$ -a.s.,

$$\nabla \cdot q_{n,\xi} = 0 \text{ and } \oint_{\square_n} q_{n,\xi} = 0.$$

II. The variational approach

A useful decomposition: for every $n \in \mathbb{N}$, up to a set of Lebesgue measure zero,

$$\square_n = \left[-\frac{3^n}{2}, \frac{3^n}{2} \right] \text{ so } \mathbb{R}^d = \coprod_{z \in 3^n \mathbb{Z}^d} (z + \square_n).$$

The finite-volume correctors: for every $n \in \mathbb{N}$,

$$-\nabla \cdot a(e_i + \nabla \phi_{n,i}) = 0 \text{ in } \square_n \text{ with } \phi_{n,i} = 0 \text{ on } \partial \square_n.$$

The finite volume flux: for every $n \in \mathbb{N}$,

$$q_{n,i} = a(e_i + \nabla \phi_{n,i}) - a(\square_n)e_i,$$

which satisfies, $\mathbb{P}$ -a.s.,

$$\nabla \cdot q_{n,i} = 0 \text{ and } \int_{\square_n} q_{n,i} = 0.$$

The finite volume flux-corrector: for 3^n -periodic skew-symmetric matrices $\sigma_{i,n}$,

$$\nabla \cdot \sigma_{n,i} = q_{n,i} \text{ in } \mathbb{R}^d,$$

with $\sigma_{n,i} = (\sigma_{n,ijk})$ for

$$-\Delta \sigma_{n,ijk} = \partial_j q_{n,ik} - \partial_k q_{n,ij} \text{ in } \mathbb{R}^d \text{ with } \int_{\square_n} \sigma_{n,ijk} = 0.$$

II. The variational approach

The equation: we have

$$-\nabla \cdot a^\varepsilon \nabla u_\varepsilon = f \text{ in } U \text{ with } u_\varepsilon = 0 \text{ on } \partial U.$$

The two-scale expansion: if $U \subseteq \square_0$ and $3^{m-1} \leq \varepsilon^{-1} < 3^m$ then $(U/\varepsilon) \subseteq \square_m$,

$$u_\varepsilon \simeq \tilde{u}_{\varepsilon,m} = u_\varepsilon - \bar{u} - \varepsilon \partial_i \phi_{m,i}^\varepsilon \partial_i \bar{u}.$$

for the finite-volume correctors

$$-\nabla \cdot a(e_i + \nabla \phi_{m,i}) = 0 \text{ in } \square_m \text{ with } \phi_{m,i} = 0 \text{ on } \partial \square_m.$$

The finite volume flux: for every $n \in \mathbb{N}$,

$$q_{n,i} = a(e_i + \nabla \phi_{n,i}) - a(\square_n) e_i,$$

which satisfies, $\mathbb{P}$ -a.s.,

$$\nabla \cdot q_{n,i} = 0 \text{ and } \oint_{\square_n} q_{n,i} = 0.$$

The finite volume flux-corrector: for 3^n -periodic skew-symmetric matrices $\sigma_{i,n}$,

$$\nabla \cdot \sigma_{n,i} = q_{n,i} \text{ in } \mathbb{R}^d,$$

with $\sigma_{n,i} = (\sigma_{n,ijk})$ for

$$-\Delta \sigma_{n,ijk} = \partial_j q_{n,ik} - \partial_k q_{n,ij} \text{ in } \mathbb{R}^d \text{ with } \oint_{\square_n} \sigma_{n,ijk} = 0.$$

II. The variational approach

The equation: we have

$$-\nabla \cdot a^\varepsilon \nabla u_\varepsilon = f \text{ in } U \text{ with } u_\varepsilon = 0 \text{ on } \partial U.$$

The two-scale expansion: the homogenization error, for $U \subseteq \square_0$ and $3^{m-1} \leq \varepsilon^{-1} < 3^m$,

$$w_\varepsilon = u_\varepsilon - \bar{u} - \varepsilon \phi_{m,i}^\varepsilon \partial_i \bar{u},$$

satisfies

$$\begin{aligned} -\nabla \cdot a^\varepsilon \nabla w_\varepsilon &= q_{m,i}^\varepsilon \cdot \nabla \partial_i \bar{u} + \nabla \cdot (a(\square_m) - \bar{a}) \nabla \bar{u} + \nabla \cdot (\varepsilon \phi_{m,i}^\varepsilon \nabla \partial_i \bar{u}) \\ &= \nabla \cdot (a(\square_m) - \bar{a}) \nabla \bar{u} + \nabla \cdot (\varepsilon \phi_{m,i}^\varepsilon - \varepsilon \sigma_{m,i}^\varepsilon) \nabla \partial_i \bar{u}. \end{aligned}$$

Sublinearity: since $3^{m-1} \leq \varepsilon^{-1} < 3^m$,

$$\begin{aligned} \|\varepsilon \phi_{m,i}^\varepsilon\|_{L^2(U)} + \|\varepsilon \sigma_{m,i}^\varepsilon\|_{L^2(U)} &= \varepsilon |U|^{\frac{1}{2}} \left(\int_{U/\varepsilon} |\phi_{m,i}|^2 + |\sigma_{m,i}|^2 \right)^{\frac{1}{2}} \\ &\leq c 3^{-m} \left(\int_{\square_m} |\phi_{m,i}|^2 + |\sigma_{m,i}|^2 \right)^{\frac{1}{2}}. \end{aligned}$$

It suffices to prove that

$$\lim_{m \rightarrow \infty} 3^{-m} \left(\|\phi_{m,i}\|_{\underline{L}^2(U)} + \|\sigma_{m,i}\|_{\underline{L}^2(U)} \right) = 0.$$

III. The multiscale Poincaré inequality

The Poincaré inequality: for every $U \subseteq \mathbb{R}^d$,

$$\|u - \langle u \rangle_U\|_{L^2(U)} \leq c \operatorname{diam}(U) \|\nabla u\|_{L^2(U)}.$$

Scale-corrected norms: we define the following volume-normalized function spaces,

$$\begin{aligned}\|u\|_{\underline{L}^2(U)} &= \left(\int_U u^2 \right)^{\frac{1}{2}}, \\ \|u\|_{\underline{H}^1(U)} &= (|U|^{-1} \int_U u^2)^{\frac{1}{2}} + \left(\int_U |\nabla u|^2 \right)^{\frac{1}{2}}, \\ \|u\|_{\underline{H}^{-1}(U)} &= \sup \left\{ \int_U uv : \|v\|_{\underline{H}^1(U)} \leq 1 \right\}.\end{aligned}$$

The scaling factor: we have that

$$\|v^\varepsilon\|_{\underline{L}^2(\varepsilon U)} = \|v\|_{\underline{L}^2(U)} \quad \text{and} \quad \|v^\varepsilon\|_{\underline{H}^1(\varepsilon U)} = \varepsilon^{-1} \|v\|_{\underline{H}^1(U)}.$$

The scaling $v \rightarrow \varepsilon v^\varepsilon$ preserves the $\underline{H}^1$ -norm.

III. The multiscale Poincaré inequality

A useful characterization of the L^2 -norm: we aim to show that

$$\|u - \langle u \rangle_U\|_{\underline{L}^2(U)} \leq c \|\nabla u\|_{\underline{H}^{-1}(U)}.$$

The Neumann solution: let w solve

$$-\Delta w = (u - \langle u \rangle_U) \text{ in } U \text{ with } \nabla w \cdot \nu = 0 \text{ on } \partial U.$$

Maximal regularity: we have that

$$\|\nabla w\|_{\underline{H}^1(U)} \leq c \|u - \langle u \rangle_U\|_{\underline{L}^2(U)},$$

using the $\nabla w \rightarrow \varepsilon(\nabla w)^\varepsilon$ scaling.

The estimate: we have that

$$\begin{aligned} \|u - \langle u \rangle_U\|_{\underline{L}^2(U)}^2 &= \int_U \nabla w \cdot \nabla u \\ &\leq \|\nabla w\|_{\underline{H}^1(U)} \|\nabla u\|_{\underline{H}^{-1}(U)} \\ &\leq c \|u - \langle u \rangle_U\|_{\underline{L}^2(U)} \|\nabla u\|_{\underline{H}^{-1}(U)}. \end{aligned}$$

For $c \in (0, \infty)$ independent of U ,

$$\|u - \langle u \rangle_U\|_{\underline{L}^2(U)}^2 \leq c \|\nabla u\|_{\underline{H}^{-1}(U)}.$$

III. The multiscale Poincaré inequality

The Lebesgue decomposition: for $m \in \mathbb{N}$ and $g \in L^1(\square_m)$,

$$g = \langle g \rangle_{\square_m} + \sum_{k=-\infty}^m \sum_{z \in 3^k \mathbb{Z}^d \cap \square_m} \sum_{z' \in 3^{k-1} \mathbb{Z}^d \cap (z + \square_k)} (\langle g \rangle_{z' + \square_{k-1}} - \langle g \rangle_{z + \square_k}) \mathbf{1}_{z' + \square_{k-1}},$$

and, therefore,

$$\int_{\square_m} fg = \langle f \rangle_{\square_m} \langle g \rangle_{\square_m} + \sum_k \overline{\sum_{z, z'}} (\langle g \rangle_{z' + \square_{k-1}} - \langle g \rangle_{z + \square_k}) \langle f \rangle_{z' + \square_{k-1}}.$$

Hölder's inequality shows that

$$\begin{aligned} \left| \int_{\square_m} fg \right| &\leq \langle f \rangle_{\square_m} \langle g \rangle_{\square_m} + \sum_k \overline{\sum_{z, z'}} |\langle g \rangle_{z' + \square_{k-1}} - \langle g \rangle_{z + \square_k}| |\langle f \rangle_{z' + \square_{k-1}}| \\ &\leq \langle f \rangle_{\square_m} \langle g \rangle_{\square_m} + \sum_k \overline{\sum_z} 3^d \int_{z + \square_k} |g - \langle g \rangle_{z + \square_k}| \overline{\sum_{z'}} |\langle f \rangle_{z' + \square_{k-1}}| \\ &\leq \langle f \rangle_{\square_m} \langle g \rangle_{\square_m} + 3^d \sum_k \left(\overline{\sum_z} \left(\int_{z + \square_k} |g - \langle g \rangle_{z + \square_k}|^2 \right)^{\frac{1}{2}} \left(\overline{\sum_{z'}} |\langle f \rangle_{z' + \square_{k-1}}|^2 \right)^{\frac{1}{2}} \right) \\ &\leq \langle f \rangle_{\square_m} \langle g \rangle_{\square_m} + 3^{d+1} \sup_{k \leq m} 3^{-k} \left(\overline{\sum_z} \left(\int_{z + \square_k} |g - \langle g \rangle_{z + \square_k}|^2 \right)^{\frac{1}{2}} \right) \\ &\quad \times \sum_k 3^k \left(\overline{\sum_{z'}} |\langle f \rangle_{z' + \square_{k-1}}|^2 \right)^{\frac{1}{2}}. \end{aligned}$$

III. The multiscale Poincaré inequality

The Poincaré inequality: since we have that, if $\|g\|_{\underline{H}^1(\square_m)} \leq 1$,

$$\begin{aligned} \sup_{k \leq m} 3^{-k} \left(\overline{\sum_z} \left(\int_{z+\square_k} |g - \langle g \rangle_{z+\square_k}| \right)^2 \right)^{\frac{1}{2}} &\leq c \sup_{k \leq m} \left(\overline{\sum_z} \int_{z+\square_k} |\nabla g|^2 \right)^{\frac{1}{2}} \\ &\leq c \|\nabla g\|_{\underline{L}^2(\square_m)} \\ &\leq c. \end{aligned}$$

The conclusion: since $\langle g \rangle_{\square_m} \leq 3^m$ if $\|g\|_{\underline{H}^1(\square_m)} \leq 1$,

$$\|f\|_{\underline{H}^{-1}(\square_m)} = \sup_{\|g\|_{\underline{H}^1(\square_m)}} \left| \int_{\square_m} fg \right| \leq \sum_{k=-\infty}^m 3^k \left(\overline{\sum_{z \in 3^k \mathbb{Z}^d \cap \square_m}} |\langle f \rangle_{z+\square_k}|^2 \right)^{\frac{1}{2}}.$$

and, using that $\|u - \langle u \rangle_U\|_{\underline{L}^2(U)}^2 \leq c \|\nabla u\|_{\underline{H}^{-1}(U)}$,

$$\|u - \langle u \rangle_{\square_m}\|_{\underline{L}^2(\square_m)}^2 \leq c \sum_{k=-\infty}^m 3^k \left(\overline{\sum_{z \in 3^k \mathbb{Z}^d \cap \square_m}} |\langle \nabla u \rangle_{z+\square_k}|^2 \right)^{\frac{1}{2}}.$$

In particular, for $c \in (0, \infty)$ independent of m ,

$$\|u - \langle u \rangle_{\square_m}\|_{\underline{L}^2(\square_m)}^2 \leq c \|\nabla u\|_{\underline{L}^2(\square_m)} + c \sum_{k=0}^m 3^k \left(\overline{\sum_{z \in 3^k \mathbb{Z}^d \cap \square_m}} |\langle \nabla u \rangle_{z+\square_k}|^2 \right)^{\frac{1}{2}}.$$

IV. Qualitative stochastic homogenization

The finite-volume correctors: for $v_{n,\xi}$ minimizing $\mu(\square_n, \xi)$,

$$-\nabla \cdot a \nabla v_{n,\xi} = 0 \text{ in } \square_n \text{ with } v_{n,\xi} = \ell_\xi \text{ on } \partial \square_n,$$

we have that $\phi_{n,\xi} = v_{n,\xi} - \ell_\xi$ and, since $\phi_{n,\xi} = 0$ on $\partial \square_n$,

$$\langle \phi_{n,\xi} \rangle_{\square_n} = 0 \text{ implies } \langle v_{n,\xi} \rangle_{\square_n} = \xi.$$

Large-scale averages: we have that, for every $n \leq m \in \mathbb{N}_0$,

$$\begin{aligned} |\langle \nabla \phi_{m,\xi} \rangle_{\square_m}|^2 &= |\langle \nabla v_{m,\xi} - \xi \rangle_{\square_m}|^2 \\ &= |\langle \nabla v_{m,\xi} \rangle_{\square_m} - \langle \nabla v_{n,\xi} \rangle_{\square_m}|^2 \\ &= |\fint_{\square_m} (\nabla v_{m,\xi} - \nabla v_{n,\xi})|^2 \\ &\leq \fint_{\square_m} |\nabla v_{m,\xi} - \nabla v_{n,\xi}|^2. \end{aligned}$$

IV. Qualitative stochastic homogenization

Large-scale averages: for the finite-volume corrector, for every $n \leq m \in \mathbb{N}_0$,

$$|\langle \nabla \phi_{m,\xi} \rangle_{\square_m}|^2 \leq \int_{\square_m} |\nabla v_{m,\xi} - \nabla v_{n,\xi}|^2.$$

Asymptotic flatness: for every $n \leq m \in \mathbb{N}_0$,

$$\begin{aligned} \frac{\lambda}{2} \int_{\square_m} |\nabla v_{m,\xi} - \nabla v_{n,\xi}|^2 &\leq \frac{1}{2} \int_{\square_m} (\nabla v_{m,\xi} - \nabla v_{n,\xi}) \cdot a(\nabla v_{m,\xi} - \nabla v_{n,\xi}) \\ &= \frac{1}{2} \int_{\square_m} \nabla v_{n,\xi} \cdot a \nabla v_{n,\xi} - \frac{1}{2} \int_{\square_m} \nabla v_{n,\xi} \cdot a \nabla v_{m,\xi} \\ &= \frac{1}{2} \int_{\square_m} \nabla v_{n,\xi} \cdot a \nabla v_{n,\xi} - \frac{1}{2} \int_{\square_m} \nabla v_{m,\xi} \cdot a \nabla v_{m,\xi} \\ &= \overline{\sum_{z \in 3^n \mathbb{Z}^d \cap \square_M}} (\mu(z + \square_n, \xi) - \mu(\square_m, \xi)) \\ &= \overline{\sum_{z \in 3^n \mathbb{Z}^d \cap \square_M}} (\xi \cdot a(z + \square_n) \xi - \xi \cdot a(\square_m) \xi) \\ &= \overline{\sum_{z \in 3^n \mathbb{Z}^d \cap \square_M}} (\xi \cdot (a(z + \square_n) - a(\square_m)) \xi). \end{aligned}$$

IV. Qualitative stochastic homogenization

Large-scale averages: for the finite-volume corrector, for every $n \leq m \in \mathbb{N}_0$,

$$|\langle \nabla \phi_{m,\xi} \rangle_{\square_m}|^2 \leq \int_{\square_m} |\nabla v_{m,\xi} - \nabla v_{n,\xi}|^2.$$

Asymptotic flatness: for every $n \leq m \in \mathbb{N}_0$,

$$\frac{\lambda}{2} \int_{\square_m} |\nabla v_{m,\xi} - \nabla v_{n,\xi}|^2 = \overline{\sum_{z \in 3^n \mathbb{Z}^d \cap \square_m}} (\xi \cdot (a(z + \square_n) - a(\square_m)) \xi).$$

The multiscale Poincaré inequality: we have that

$$\begin{aligned} & \|\phi_{m,\xi} - \langle \phi_{m,\xi} \rangle_{\square_m}\|_{\underline{L}^2(\square_m)} \\ & \leq c \left(\|\nabla \phi_{m,\xi}\|_{\underline{L}^2} + \sum_{k=0}^m 3^k \left(\overline{\sum_{z \in 3^k \mathbb{Z}^d \cap \square_m}} |\langle \nabla \phi_{m,\xi} \rangle_{z + \square_k}|^2 \right)^{\frac{1}{2}} \right) \\ & \leq c \left(|\xi|^2 + \sum_{k=0}^{m-1} 3^k \left(\overline{\sum_{z \in 3^n \mathbb{Z}^d \cap \square_m}} (\xi \cdot (a(z + \square_n) - a(\square_m)) \xi) \right)^{\frac{1}{2}} \right) \\ & \leq c |\xi|^2 \left(1 + \sum_{k=0}^{m-1} 3^k \left(|a(\square_m) - \overline{\sum_{z \in 3^n \mathbb{Z}^d \cap \square_m}} a(z + \square_n)| \right)^{\frac{1}{2}} \right) \\ & \leq c |\xi|^2 \left(1 + \sum_{k=0}^m 3^k \left(\left| \overline{\sum_{z \in 3^n \mathbb{Z}^d \cap \square_m}} a(z + \square_n) - \bar{a} \right| \right)^{\frac{1}{2}} \right). \end{aligned}$$

IV. Qualitative stochastic homogenization

Large-scale averages: for the finite-volume corrector, for every $n \leq m \in \mathbb{N}_0$,

$$|\langle \nabla \phi_{m,\xi} \rangle_{\square_m}|^2 \leq \int_{\square_m} |\nabla v_{m,\xi} - \nabla v_{n,\xi}|^2.$$

Asymptotic flatness: for every $n \leq m \in \mathbb{N}_0$,

$$\frac{\lambda}{2} \int_{\square_m} |\nabla v_{m,\xi} - \nabla v_{n,\xi}|^2 = \sum_{z \in 3^n \mathbb{Z}^d \cap \square_m} (\xi \cdot (a(z + \square_n) - a(\square_m)) \xi).$$

The multiscale Poincaré inequality: we have that

$$\begin{aligned} & \|\phi_{m,\xi} - \langle \phi_{m,\xi} \rangle_{\square_m}\|_{\underline{L}^2(\square_m)} \\ & \leq c|\xi|^2 \left(1 + \sum_{k=0}^m 3^k \left(\left| \sum_{z \in 3^n \mathbb{Z}^d \cap \square_m} a(z + \square_n) - \bar{a} \right| \right)^{\frac{1}{2}} \right). \end{aligned}$$

Sublinearity: by the subadditive ergodic theorem and dominated convergence,

$$\begin{aligned} & \limsup_{m \rightarrow \infty} 3^{-m} \left(\|\phi_{m,\xi} - \langle \phi_{m,\xi} \rangle_{\square_m}\|_{\underline{L}^2(\square_m)} \right) \\ & \leq \limsup_{m \rightarrow \infty} c|\xi|^2 \left(3^{-m} + \sum_{k=0}^m 3^{k-m} \left(\left| \sum_{z \in 3^n \mathbb{Z}^d \cap \square_m} a(z + \square_n) - \bar{a} \right| \right)^{\frac{1}{2}} \right) \\ & = 0. \end{aligned}$$

IV. Qualitative stochastic homogenization

The finite volume flux: for every $n \in \mathbb{N}$,

$$q_{m,i} = a(e_i + \nabla \phi_{m,i}) - a(\square_m)e_i.$$

The finite volume flux-corrector: for 3^m -periodic skew-symmetric matrices $\sigma_{i,m}$ with

$$-\Delta \sigma_{m,ijk} = \partial_j q_{m,ik} - \partial_k q_{m,ij} \text{ in } \mathbb{R}^d \text{ with } \oint_{\square_n} \sigma_{m,ijk} = 0.$$

The L^2 -norm: for the 3^m -periodic function w satisfying

$$-\Delta w = \sigma_{m,ijk} \text{ in } \mathbb{R}^d \text{ with } \langle w \rangle_{\square_m} = 0,$$

we have using the maximal regularity estimate that

$$\begin{aligned} \|\sigma_{m,ijk}\|_{\underline{L}^2(\square_m)}^2 &\leq \oint_{\square_m} \nabla w \cdot \nabla \sigma_{m,ijk} = \oint_{\square_m} q_{m,ij} \partial_k w - q_{m,ik} \partial_j w \\ &\leq \|q_{m,i}\|_{\underline{H}^{-1}(\square_m)} \|\nabla w\|_{\underline{H}^1(\square_m)} \\ &\leq \|q_{m,i}\|_{\underline{H}^{-1}(\square_m)} \|\sigma_{m,ijk}\|_{\underline{L}^2(\square_m)}, \end{aligned}$$

and, therefore, $\|\sigma_{m,ijk}\|_{\underline{L}^2(\square_m)} \leq \|q_{m,i}\|_{\underline{H}^{-1}(\square_m)}$.

Sublinearity: applying the multiscale Poincaré inequality to $q_{m,i}$.

$$\limsup_{m \rightarrow \infty} 3^{-m} \left(\|\sigma_{m,\xi} - \langle \sigma_{m,\xi} \rangle_{\square_m}\|_{\underline{L}^2(\square_m)} \right) = 0.$$

IV. Qualitative stochastic homogenization

Qualitative stochastic homogenization

Let $U \subseteq \square_0$ be a smooth, bounded domain, let $f \in C^\alpha(\overline{U})$ for some $\alpha \in (0, 1)$, and let $a: \mathbb{R}^d \times \Omega \rightarrow \mathbb{R}^{d \times d}$ be a symmetric, uniformly elliptic, stationary, and ergodic diffusion matrix. Then, there exists a deterministic, uniformly elliptic $\bar{a} \in \mathbb{R}^{d \times d}$ defined $\mathbb{P}$ -a.s. by

$$\frac{1}{2}(\xi \cdot \bar{a}\xi) = \lim_{m \rightarrow \infty} \mu(\square_m, \xi),$$

such that, $\mathbb{P}$ -a.s., the local two scale expansion, for $3^{m-1} \leq \varepsilon^{-1} < 3^m$,

$$w_\varepsilon = u_\varepsilon - \varepsilon \phi_{m,i}^\varepsilon \partial_i \bar{u},$$

defined by the finite-volume correctors $\phi_{m,i}$ and macroscopic limit

$$-\nabla \cdot \bar{a} \nabla \bar{u} = f \text{ in } U \text{ with } \bar{u} = 0 \text{ on } \partial U,$$

satisfies

$$\lim_{\varepsilon \rightarrow 0} \int_U |\nabla w_\varepsilon|^2 = 0 \text{ and, so, } \nabla u_\varepsilon \simeq \nabla \bar{u} + \partial_i \bar{u} \nabla \phi_i^\varepsilon.$$

Polarization: for every $\xi, \xi' \in \mathbb{R}^d$,

$$\xi \cdot \bar{a} \xi' = \frac{1}{4} \left((\xi + \xi') \cdot \bar{a} (\xi + \xi') - (\xi - \xi') \cdot \bar{a} (\xi - \xi') \right).$$

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