

Diffusions in random environment and stochastic homogenization

Benjamin Fehrman

University of Oxford

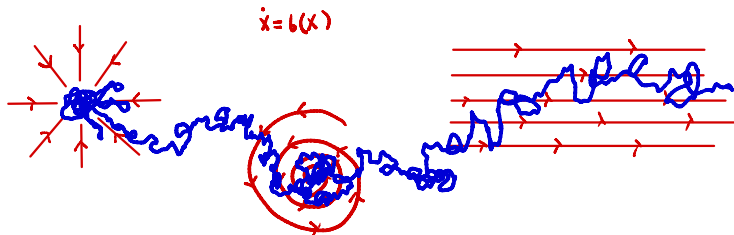
20-24 September 2021

I. Diffusion processes and PDEs

In these talks, we will consider the longtime behavior of a diffusion process

$$dX_t = \sigma(X_t) dB_t + b(X_t) dt \text{ for } t \in (0, \infty).$$

- σ quantifies the diffusion
 - thermal fluctuations / microscopic collisions driving a Brownian particle
- b quantifies the drift
 - mean macroscopic motion / wind or current in a fluid flow



I. Diffusion processes and PDEs

If X_t solves

$$dX_t = \sigma(X_t) dB_t + b(X_t) dt \text{ for } t \in (0, \infty),$$

the central limit scaling $X^\varepsilon = \varepsilon X_{t/\varepsilon^2}$ solves

$$dX_t^\varepsilon = \sigma(X_t^\varepsilon/\varepsilon) dW_t^\varepsilon + \varepsilon^{-1} b(X_t^\varepsilon/\varepsilon) dt.$$

What happens (in law) as $\varepsilon \rightarrow 0$?

- (*diffusive*) If $\sigma = I$ and $b = 0$ then, in law for every $\varepsilon \in (0, 1)$,

$$X_t^\varepsilon = B_t.$$

- (*ballistic*) If $\sigma = I$ and $b = \bar{b} \in \mathbb{R}^d \setminus \{0\}$ then, in law for every $\varepsilon \in (0, 1)$,

$$X_t^\varepsilon = B_t + \varepsilon^{-1} t \bar{b} \text{ and almost surely } |X_t^\varepsilon| \rightarrow \infty \text{ as } \varepsilon \rightarrow 0.$$

- (*degenerate / trapped*) For the Ornstein-Uhlenbeck process

$$dX_t = dB_t - X_t dt \text{ and } dX_t^\varepsilon = dW_t^\varepsilon - \varepsilon^{-2} X_t^\varepsilon dt,$$

and

$$|X_t^\varepsilon| \rightarrow 0 \text{ almost surely as } \varepsilon \rightarrow 0.$$

I. Diffusion processes and PDEs

We are interested in the behavior of X_t in *law*.

If X_t solves

$$dX_t = \sigma(X_t) dB_t + b(X_t) dt \text{ for } t \in (0, \infty),$$

how can we characterize

$$\mathbb{P}[X_t \in A] \text{ for every measurable } A \subseteq \mathbb{R}^d?$$

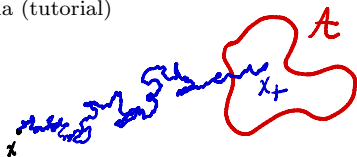
The *Feynman-Kac Formula*: if ρ solves the equation

$$\partial_t \rho = \text{tr}(a \nabla^2 \rho) + b \cdot \nabla \rho \text{ in } \mathbb{R}^d \times (0, \infty) \text{ with } \rho(\cdot, 0) = \rho_0,$$

for the covariance matrix $a = \frac{1}{2} \sigma \sigma^2$, then we have the formula

$$\rho(x, t) = \mathbb{E}_x [\rho_0(X_t)].$$

- the heat equation and Brownian motion
- the solution is the average of the initial data with respect to the diffusion
 - regularizing / smoothing properties of parabolic equations
- proof using Itô's formula (tutorial)



I. Diffusion processes and PDEs

In the central limit scaling $X_t^\varepsilon = \varepsilon X_{t/\varepsilon^2}$,

$$dX_t^\varepsilon = \sigma(X_t^\varepsilon/\varepsilon) dW_t^\varepsilon + \varepsilon^{-1} b(X_t^\varepsilon/\varepsilon) dt,$$

and the solution ρ^ε of the equation

$$\partial_t \rho^\varepsilon = \text{tr}(a(x/\varepsilon) \nabla^2 \rho^\varepsilon) + \varepsilon^{-1} b(x/\varepsilon) \cdot \nabla \rho^\varepsilon \quad \text{in } \mathbb{R}^d \times (0, \infty) \quad \text{with } \rho(\cdot, 0) = \rho_0,$$

for $a = \frac{1}{2} \sigma \sigma^2$ satisfies

$$\rho(x, t) = \mathbb{E}_x [\rho_0(X_t^\varepsilon)] = \mathbb{E}_{\frac{x}{\varepsilon}} [\rho_0(\varepsilon X_{t/\varepsilon^2})].$$

- (*diffusive*) If $\sigma = I$ and $b = 0$ then, for every $\varepsilon \in (0, 1)$,

$$\rho^\varepsilon = \bar{\rho} \quad \text{for } \partial_t \bar{\rho} = \frac{1}{2} \Delta \bar{\rho}.$$

- (*ballistic*) If $\sigma = I$ and $b = \bar{b} \in \mathbb{R}^d \setminus \{0\}$ then

$$\left(\lim_{\varepsilon \rightarrow 0} \rho^\varepsilon(x, t) \right) = \left(\lim_{s \rightarrow \infty} \rho_0(x + s\bar{b}) \right).$$

- (*degenerate / trapped*) In the case of the Ornstein-Uhlenbeck process,

$$\partial_t \rho^\varepsilon = \frac{1}{2} \Delta \rho^\varepsilon - \varepsilon^{-2} x \cdot \nabla \rho^\varepsilon,$$

and $(\lim_{\varepsilon \rightarrow 0} \rho^\varepsilon(x, t)) = \rho_0(0)$.

I. Diffusion processes and PDEs

Characterizing the limiting behavior, as $\varepsilon \rightarrow 0$, of the solution

$$dX_t^\varepsilon = \sigma(X_t^\varepsilon/\varepsilon) dW_t^\varepsilon + \varepsilon^{-1} b(X_t^\varepsilon/\varepsilon) dt,$$

in law is equivalent to characterizing the limiting behavior, as $\varepsilon \rightarrow 0$, of the solutions

$$\partial_t \rho^\varepsilon = \text{tr}(a(x/\varepsilon) \nabla^2 \rho^\varepsilon) + \varepsilon^{-1} b(x/\varepsilon) \cdot \nabla \rho^\varepsilon,$$

for arbitrary smooth initial data.

- The Feynman-Kac formula:

$$\rho^\varepsilon(x, t) = \mathbb{E}_x [\rho_0(X_t^\varepsilon)].$$

- As $\varepsilon \rightarrow 0$, we have $X^\varepsilon \rightarrow \bar{X}$ in law, for \bar{X} solving

$$d\bar{X}_t = \bar{\sigma} dB_t \text{ for some } \bar{\sigma} \in \mathbb{R}^{d \times d},$$

if and only if we have $\rho^\varepsilon \rightarrow \bar{\rho}$, for $\bar{\rho}$ solving

$$\partial_t \bar{\rho} = \text{tr}(\bar{a} \nabla^2 \bar{\rho}) \text{ for } \bar{a} = \frac{1}{2} \bar{\sigma} \bar{\sigma}^t.$$

- The divergence-form case / a reversible diffusion:

$$-\nabla \cdot (a(x/\varepsilon) \nabla \rho^\varepsilon) = -\text{tr}(a(x/\varepsilon) \nabla^2 \rho^\varepsilon) - \varepsilon^{-1} (\nabla \cdot a^t(x/\varepsilon)) \cdot \nabla \rho^\varepsilon.$$

II. Ergodic properties of diffusions on the torus

We will restrict (for now) to periodic coefficient fields.

- For 1-periodic coefficients σ and b , we have the diffusion X on \mathbb{R}^d :

$$dX_t = \sigma(X_t) dB_t + b(X_t) dt \text{ in } \mathbb{R}^d.$$

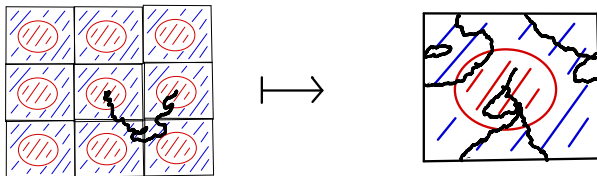
Lift this to a diffusion \bar{X} on the torus \mathbb{T}^d :

$$d\bar{X}_t = \sigma(\bar{X}_t) dB_t + b(\bar{X}_t) dt \text{ in } \mathbb{T}^d.$$

- For $\rho_0 \in C^\infty(\mathbb{T}^d)$, the function

$$\rho(x, t) = (\bar{P}_t \rho_0)(x) = \mathbb{E}_x[\rho_0(\bar{X}_t)] \text{ solves } \partial_t \rho = \text{tr}(a \nabla^2 \rho) + b \cdot \nabla \rho \text{ in } \mathbb{T}^d.$$

- What are the averaging / ergodic properties of the semigroup \bar{P}_t ?



II. Ergodic properties of diffusions on the torus

The invariant measure [Section 3.2, Asym. Anal. for Per. Struct.]

Assume that σ and b are sufficiently regular, and assume that $a = \frac{1}{2}\sigma\sigma^t$ is uniformly elliptic: there exist $\lambda \leq \Lambda \in (0, \infty)$ such that

$$\lambda |\xi|^2 \leq \langle a(x)\xi, \xi \rangle \leq \Lambda |\xi|^2 \quad \text{for every } x \in \mathbb{T}^d \text{ and } \xi \in \mathbb{R}^d.$$

Then there exists a unique probability measure π on \mathbb{T}^d and constants $c, \rho \in (0, \infty)$ such that, for every $f \in L^\infty(\mathbb{T}^d)$,

$$\sup_{x \in \mathbb{T}^d} \left| \mathbb{E}_x [f(\bar{X}_t)] - \int_{\mathbb{T}^d} f(y) \pi(dy) \right| \leq c \|f\|_{L^\infty(\mathbb{T}^d)} \exp(-\rho t).$$

- *uniform ellipticity* yields exponential convergence to the invariant distribution
- The semigroup \bar{P}_t on functions defines an adjoint semigroup on \bar{P}_t^* on measures:

$$\int_{\mathbb{T}^d} f(y) (\bar{P}_t^* \mu)(dy) := \int_{\mathbb{T}^d} \bar{P}_t f(y) \mu(dy) = \int_{\mathbb{T}^d} \mathbb{E}_y [f(\bar{X}_t)] \mu(dy).$$

- *Invariance*: we have that $(\bar{P}_t^* \pi) = \pi$ for every $t \in [0, \infty)$, since

$$\int_{\mathbb{T}^d} \bar{P}_t f(y) \pi(dy) = \int_{\mathbb{T}^d} f(y) \pi(dy) \quad \text{for every } t \in [0, \infty),$$

- *Uniqueness / absolute continuity with respect to Lebesgue measure* (tutorial)

II. Ergodic properties of diffusions on the torus

For 1-periodic coefficients,

$$d\bar{X}_t = \sigma(\bar{X}_t) dB_t + b(\bar{X}_t) dt \quad \text{in } \mathbb{T}^d.$$

- We have the unique, mutually, absolutely continuous invariant measure $(\bar{P}_t^* \pi) = \pi$:

$$\int_{\mathbb{T}^d} E_y [F(\bar{X}_t)] \pi(dy) = \int_{\mathbb{T}^d} \bar{P}_t f(y) \pi(dy) = \int_{\mathbb{T}^d} f(y) \pi(y) dy.$$

- By absolute continuity, the invariant measure π has a positive density m in $L^1(\mathbb{T}^d)$:

$$d\pi = m(y) dy.$$

- By Feynman-Kac if $\partial_t \rho = \text{tr}(a \nabla^2 \rho) + b \cdot \nabla \rho$ in \mathbb{T}^d then

$$\int_{\mathbb{T}^d} \rho_0(y) m(y) dy = \int_{\mathbb{T}^d} \rho(y, t) m(y) dy \quad \text{for every } t \in [0, \infty).$$

- For the differential operator

$$\mathcal{L}g = \text{tr}(a \nabla^2 g) + b \cdot \nabla g \quad \text{and its adjoint } \mathcal{L}^*g = (a_{ij}g)_{x_i x_j} - \nabla \cdot (gb),$$

we have that

$$0 = \partial_t \left(\int_{\mathbb{T}^d} \rho(y, t) m(y) dy \right) = \int_{\mathbb{T}^d} (\mathcal{L}\rho(y, t)) m(y) dy = \int_{\mathbb{T}^d} \rho(x, t) \mathcal{L}^* m(y) dy.$$

- The density solves the adjoint equation $\mathcal{L}^* m = 0$.

II. Ergodic properties of diffusions on the torus

The Fredholm alternative [Section 3.3, Asym. Anal. for Per. Struct.]

Let σ and b be sufficiently regular, and let a be uniformly elliptic. Consider the equations

$$\operatorname{tr}(a\nabla^2\rho) + b \cdot \nabla\rho = \mathcal{L}\rho = 0 \text{ in } \mathbb{T}^d, \quad (1)$$

and

$$(a_{ij}z)_{x_i x_j} - \nabla \cdot (zb) = \mathcal{L}^*z = 0 \text{ in } \mathbb{T}^d. \quad (2)$$

Then up to a multiplicative constant there exists a unique solution of (1) and (2) (namely, $\rho = 1$ and $z = m$, the density invariant measure). Furthermore, for $\phi, \psi \in L^\infty(\mathbb{T}^d)$ satisfying

$$\int_{\mathbb{T}^d} \phi(y)m(y) dy = 0 \text{ and } \int_{\mathbb{T}^d} \psi(y) dy = 0,$$

there exist unique solutions to the equations

$$\mathcal{L}z = \phi \text{ in } \mathbb{T}^d \text{ with } \int_{\mathbb{T}^d} z dy = 0 \text{ and } \mathcal{L}^*w = \psi \text{ in } \mathbb{T}^d \text{ with } \int_{\mathbb{T}^d} w(y) dy = 1.$$

- We can solve $\mathcal{L}z = \phi$ provided ϕ is *orthogonal* to the kernel of \mathcal{L}^* .

— Orthogonality is *necessary*: if $\mathcal{L}z = \phi$ and $\mathcal{L}^*m = 0$ then

$$\langle \phi, m \rangle_{L^2(\mathbb{T}^d)} = \langle \mathcal{L}z, m \rangle_{L^2(\mathbb{T}^d)} = \langle z, \mathcal{L}^*m \rangle_{L^2(\mathbb{T}^d)} = 0.$$

— *Sufficiency* relies strongly on *compactness*.

II. Ergodic properties of diffusions on the torus

Examples of invariant measures m :

- Divergence-form equations / reversible diffusions:

$$\mathcal{L}^* m = -\nabla \cdot a^t \nabla m = 0 \text{ implies that } m = 1.$$

- Pure diffusions in one-dimension: the case $b = 0$ and $d = 1$,

$$\mathcal{L}^* m = (am)_{xx} = 0 \text{ implies that } m = \langle a^{-1} \rangle^{-1} \frac{1}{a}$$

$$\text{for } \langle a^{-1} \rangle = \int_{\mathbb{T}^1} a^{-1}.$$

- In general, for higher dimensions and nonzero drift, they are complicated.

Using the Green's function representation and Fubini's theorem,

$$\begin{aligned} \int_{\mathbb{T}^d} \mathbb{E}_y [f(X_t)] m(y) dy &= \int_{\mathbb{T}^d} \int_{\mathbb{T}^d} \bar{p}_t(y, x) f(x) m(y) dx dy \\ &= \int_{\mathbb{T}^d} f(x) \left(\int_{\mathbb{T}^d} \bar{p}_t(y, x) m(y) dy \right) dx = \int_{\mathbb{T}^d} f(x) m(x) dx. \end{aligned}$$

We have that, for every $t \in [0, \infty)$ and $x \in \mathbb{T}^d$,

$$m(x) = \int_{\mathbb{T}^d} \bar{p}_t(y, x) m(y) dy.$$

III. Homogenization of pure diffusions

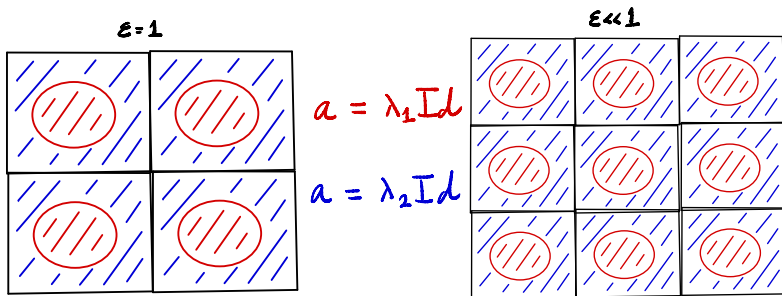
Consider the pure diffusion $dX_t = \sigma(X_t) dB_t$ and the central limit scaling

$$dX_t^\varepsilon = \sigma(X_t/\varepsilon) dW_t^\varepsilon,$$

and the corresponding equation

$$\partial_t \rho^\varepsilon = \text{tr}(a(x/\varepsilon) \nabla^2 \rho^\varepsilon),$$

for $a = \frac{1}{2} \sigma \sigma^2$.



What happens, for example, if $\lambda_1 \rightarrow 0$ and $\lambda_2 \rightarrow \infty$?

III. Homogenization of pure diffusions

Consider the pure diffusion $dX_t = \sigma(X_t) dB_t$ and the central limit scaling

$$dX_t^\varepsilon = \sigma(X_t/\varepsilon) dW_t^\varepsilon$$

and the corresponding equation

$$\partial_t \rho^\varepsilon = \text{tr}(a(x/\varepsilon) \nabla^2 \rho^\varepsilon)$$

for $a = \frac{1}{2} \sigma \sigma^2$.

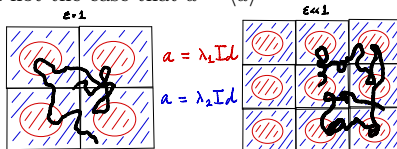
- *Homogenization*: identify $\bar{a} \in \mathbb{R}^{d \times d}$ such that

$$\rho^\varepsilon \rightarrow \bar{\rho}$$

for $\bar{\rho}$ the solution of

$$\partial_t \bar{\rho} = \text{tr}(\bar{a} \nabla^2 \bar{\rho}).$$

- Equivalently, in law, $X_t^\varepsilon \rightarrow \bar{X}$ for $d\bar{X}_t = \bar{\sigma} dB_t$.
- a complicated, nonlinear averaging
 - what is \bar{a} ?
 - it is very much not the case that $\bar{a} = \langle a \rangle$



- We have $\langle a \rangle \rightarrow \frac{1}{2} \lambda_2 I$ as $\lambda_1 \rightarrow 0$, while $\bar{a} \rightarrow 0$.

III. Homogenization of pure diffusions

The asymptotic expansion:

$$\tilde{\rho}^\varepsilon(x, t) = \bar{\rho}(x, t) + \varepsilon \rho_1(x, x/\varepsilon, t) + \varepsilon^2 \rho_2(x, x/\varepsilon, t) + \dots$$

- Evaluating the equation, keeping terms of order ε^{-1} , ε^0 , and ε ,

$$\begin{aligned} \partial_t \tilde{\rho}^\varepsilon - \operatorname{tr}(a(x/\varepsilon) \nabla^2 \tilde{\rho}^\varepsilon) &= \varepsilon^{-1} \operatorname{tr}(a(x/\varepsilon) \nabla_y^2 \rho_1) \\ &\quad \partial_t \bar{\rho} - \operatorname{tr}(a(x/\varepsilon) (\nabla_x^2 \bar{\rho} + \nabla_{xy}^2 \rho_1 + \nabla_y^2 \rho_2)) \\ &\quad + \varepsilon \partial_t \rho_1 - \varepsilon \operatorname{tr}(a(x/\varepsilon) (\nabla_x^2 \rho_1 + \nabla_{xy}^2 \rho_2)). \end{aligned}$$

- We conclude that $\rho_1 = 0$, which is very much related to the fact that

$$X_t = \sigma(X_t) dB_t \text{ is a martingale,}$$

and therefore have that

$$\partial_t \bar{\rho} = \operatorname{tr}(a(x/\varepsilon) (\nabla_x^2 \bar{\rho} + \nabla_y^2 \rho_2)) + O(\varepsilon).$$

- *Separation of scales*: we make the ansatz that $\rho_2(x, y, t) = \sum_{i,j=1}^d w_{ij}(y) \partial_{ij}^2 \bar{\rho}$ so that

$$\partial_t \bar{\rho} = \operatorname{tr}(a(x/\varepsilon) (e_{ij} + \nabla^2 w_{ij}(x/\varepsilon))) \partial_{ij}^2 \bar{\rho} + O(\varepsilon).$$

- Solvability / Fredholm alternative requires that

$$\operatorname{tr}(a(y) (e_{ij} + \nabla^2 w_{ij}(y))) = \langle a_{ij}, m \rangle_{L^2(\mathbb{T}^d)} \text{ in } \mathbb{T}^d,$$

and

$$\partial_t \bar{\rho} = \operatorname{tr}(\bar{a} \nabla^2 \bar{\rho}) \text{ for } \bar{a} = \int_{\mathbb{T}^d} a(y) m(y) dy.$$

III. Homogenization of pure diffusions

The asymptotic expansion:

$$\tilde{\rho}^\varepsilon(x, t) = \bar{\rho}(x, t) + \varepsilon^2 w_{ij}(x/\varepsilon) \nabla^2 \bar{\rho}(x, t) + \dots$$

- We define the second-order correctors:

$$\operatorname{tr}(a(y)(e_{ij} + \nabla^2 w_{ij}(y))) = \langle a_{ij}, m \rangle_{L^2(\mathbb{T}^d)} \quad \text{in } \mathbb{T}^d.$$

- We define the homogenized solution

$$\partial_t \bar{\rho} = \operatorname{tr}(\bar{a} \nabla^2 \bar{\rho}) \quad \text{for } \bar{a} = \int_{\mathbb{T}^d} a(y) m(y) dy.$$

- The asymptotic expansion $\tilde{\rho}^\varepsilon$ satisfies

$$\partial_t \tilde{\rho}^\varepsilon = \operatorname{tr}(a(x/\varepsilon) \nabla^2 \tilde{\rho}^\varepsilon) + O(\varepsilon) \quad \text{in } \mathbb{R}^d \times (0, \infty) \quad \text{with } \tilde{\rho}^\varepsilon = \rho_0 + O(\varepsilon^2).$$

- The difference $z^\varepsilon = \rho^\varepsilon - \tilde{\rho}^\varepsilon$ solves

$$\partial_t z^\varepsilon = \operatorname{tr}(a(x/\varepsilon) \nabla^2 z^\varepsilon) + O(\varepsilon) \quad \text{in } \mathbb{R}^d \times (0, \infty) \quad \text{with } z^\varepsilon = O(\varepsilon^2).$$

- Since $\tilde{\rho}^\varepsilon = \bar{\rho}(x, t) + O(\varepsilon^2)$, the comparison principle proves that, as $\varepsilon \rightarrow 0$,

$$\rho^\varepsilon \rightarrow \bar{\rho} \quad \text{for } \partial_t \bar{\rho} = \operatorname{tr}(\bar{a} \nabla^2 \bar{\rho}).$$

Or, equivalently, that in law the processes X_t^ε converges in law to $\bar{\sigma} B_t$ for $\bar{a} = \frac{1}{2} \bar{\sigma} \bar{\sigma}^2$.

III. Homogenization of pure diffusions

Homogenization of pure diffusions

Let b and σ be sufficiently regular, and let a be uniformly elliptic. Then, for every $\rho_0 \in C_c^\infty(\mathbb{R}^d)$, the solutions

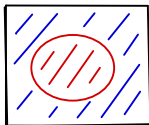
$$\partial_t \rho^\varepsilon = \operatorname{tr}(a(x/\varepsilon) \nabla^2 \rho^\varepsilon) \quad \text{in } \mathbb{R}^d \times (0, \infty) \quad \text{with } \rho^\varepsilon(\cdot, 0) = \rho_0,$$

converge, as $\varepsilon \rightarrow 0$, to the solution

$$\partial_t \bar{\rho} = \operatorname{tr}(\bar{a} \nabla^2 \bar{\rho}) \quad \text{in } \mathbb{R}^d \times (0, \infty) \quad \text{with } \bar{\rho}(\cdot, 0) = \rho_0,$$

for $\bar{a} = \int_{\mathbb{T}^d} a(y) m(y) dy$.

- The *homogenized / effective* matrix is the average of the original matrix with respect to the invariant measure.
 - uniformly elliptic with ellipticity constants determined by m
 - emphasizes regions of *small diffusion / traps / degeneracies*
- In one-dimension, we have that $m(y) = \langle a^{-1} \rangle^{-1} a^{-1}$ and $\bar{a} = \langle a^{-1} \rangle^{-1}$.



$$a = \lambda_1 \operatorname{Id}$$

$$a = \lambda_2 \operatorname{Id}$$

$$\text{If } 0 < \lambda_1 \ll \lambda_2,$$

$$m \gg 1$$

$$0 < m \ll 1.$$

IV. Periodic homogenization of divergence form equations

A random uniformly elliptic, 1-periodic coefficient field $a: \mathbb{T}^d \rightarrow \mathbb{R}^{d \times d}$.

- *uniform ellipticity*: there exists $\lambda \leq \Lambda \in (0, \infty)$ such that

$$\lambda |\xi|^2 \leq \langle a(x)\xi, \xi \rangle \leq \Lambda |\xi|^2 \quad \text{for every } x \in \mathbb{T}^d \text{ and } \xi \in \mathbb{R}^d.$$

- The solution ρ^ε of

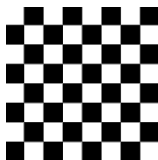
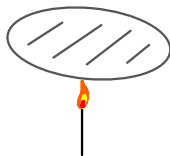
$$-\nabla \cdot a(x/\varepsilon)\nabla \rho^\varepsilon = f \text{ in } U \text{ with } \rho^\varepsilon = g \text{ on } \partial U.$$

describes the evolution of a system satisfying

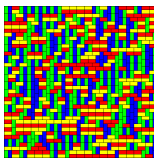
$$\oint_{B_r(x)} a(y/\varepsilon, \omega) \nabla \rho^\varepsilon(y) \cdot \nu = \int_{B_r(x)} f(y).$$

- For symmetric a we have the variational formulation and Feynman-Kac formula

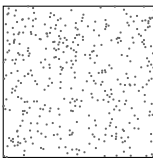
$$\inf_{v \in g + H_0^1(U)} \left(\int_U \langle a(x/\varepsilon) \nabla u, \nabla u \rangle - f u \, dx \right) \quad \text{and} \quad \rho^\varepsilon(x) = \mathbb{E}_x \left[g(X_{\tau_U}^\varepsilon) + \int_0^{\tau_U} f(X_s^\varepsilon) \, ds \right].$$



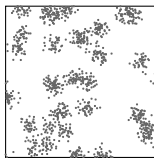
Periodic



Random tile



Poisson cloud



Cluster

IV. Periodic homogenization of divergence form equations

A *weak solution* of the equation

$$-\nabla \cdot a(x/\varepsilon)\nabla\rho^\varepsilon = f \text{ in } U \text{ with } \rho^\varepsilon = g \text{ on } \partial U.$$

is a function $\rho^\varepsilon \in H^1(U)$ that satisfies, for every $\psi \in C_c^\infty(U)$,

$$\int_U a(x/\varepsilon)\nabla\rho^\varepsilon \cdot \nabla\psi \, dx = \int_U \psi f \text{ with } \rho^\varepsilon = g \text{ on } \partial U.$$

Consider the asymptotic expansion

$$\tilde{\rho}^\varepsilon(x) = \bar{\rho}(x) + \varepsilon\rho_1(x, x/\varepsilon) + \varepsilon^2\rho_2(x, x/\varepsilon) + \dots$$

Ignoring terms of order ε and smaller, and evaluating the equation on $\tilde{\rho}^\varepsilon$,

$$\int_U a(x/\varepsilon)(\nabla_x\bar{\rho} + \nabla_y\rho_1) \cdot \nabla\psi \, dx = \int_U \psi f.$$

- Identify the equation satisfied by $\bar{\rho}$.
- Separation of scales:

$$\rho_1(x, x/\varepsilon) = \phi_i(x/\varepsilon)\partial_i\bar{\rho}(x).$$

IV. Periodic homogenization of divergence form equations

Defined by *first-order correctors* ϕ_i , the asymptotic expansion

$$\tilde{\rho}^\varepsilon = \bar{\rho} + \varepsilon \phi_i(x/\varepsilon) \partial_i \bar{\rho}(x) + \dots$$

satisfies, up to terms of order ε ,

$$\begin{aligned} \int_U a(x/\varepsilon) (\nabla_x \bar{\rho} + \nabla_y \rho_1) \cdot \nabla \psi \, dx &= \int_U a(x/\varepsilon) ((e_i + \nabla \phi_i(x/\varepsilon)) \partial_i \bar{\rho}) \cdot \nabla \psi \, dx \\ &= - \int_U \psi a(x/\varepsilon) (e_i + \nabla \phi_i(x/\varepsilon)) \cdot \nabla \partial_i \bar{\rho} = \int_U \psi f. \end{aligned}$$

As $\varepsilon \rightarrow 0$,

$$a(x/\varepsilon) (e_i + \nabla \phi_i(x/\varepsilon)) \rightharpoonup \langle a(e_i + \nabla \phi_i) \rangle =: \bar{a} e_i \text{ weakly in } L^2(\mathbb{T}^d),$$

and formally we have that

$$-\nabla \cdot a(y) (e_i + \nabla \phi_i(y)) = 0 \text{ in } \mathbb{T}^d,$$

and that $\bar{\rho}$ solves

$$-\nabla \cdot \bar{a} \nabla \bar{\rho} = f \text{ in } U \text{ with } \bar{\rho} = g \text{ on } \partial U.$$

- The functions $x_i + \phi_i(x)$ are a -harmonic
 - A *first order Liouville theorem*: any subquadratic solution of

$$-\nabla \cdot a \nabla \rho = 0 \text{ on } \mathbb{R}^d \text{ satisfies } \rho(x) = c + \xi \cdot x + \phi_\xi(x).$$

IV. Periodic homogenization of divergence form equations

The perturbed test function method [See Section 3, notes]

Let $a: \mathbb{T}^d \rightarrow \mathbb{R}^{d \times d}$ be bounded, measurable, and uniformly elliptic and let $\rho_0 \in L^2(\mathbb{R}^d)$. For every $\varepsilon \in (0, 1)$ let ρ^ε be the unique weak solution of

$$-\nabla \cdot a(x/\varepsilon)\nabla \rho^\varepsilon = f \text{ in } U \text{ with } \rho^\varepsilon = g \text{ on } \partial U.$$

Then, as $\varepsilon \rightarrow 0$,

$$\rho^\varepsilon \rightharpoonup \bar{\rho} \text{ weakly in } H^1(\mathbb{R}^d),$$

for $\bar{\rho}$ satisfying

$$-\nabla \cdot \bar{a}\nabla \bar{\rho} = f \text{ in } U \text{ with } \bar{\rho} = g \text{ on } \partial U,$$

for the effective coefficient field $\bar{a} \in \mathbb{R}^{d \times d}$ defined by

$$\bar{a}e_i = \langle a(e_i + \nabla \phi_i) \rangle \text{ for } -\nabla \cdot a(e_i + \nabla \phi_i) = 0 \text{ in } \mathbb{T}^d.$$

- compensated compactness methods
 - the div-curl lemma
- existence / uniqueness of ϕ_i follows by Fredholm or the Lax-Milgram lemma

IV. Periodic homogenization of divergence form equations

The solution ρ^ε of

$$-\nabla \cdot a(x/\varepsilon)\nabla \rho^\varepsilon = f \text{ in } U \text{ with } \rho^\varepsilon = g \text{ on } \partial U.$$

We have formally that

$$\rho^\varepsilon \simeq \tilde{\rho}^\varepsilon = \bar{\rho} + \varepsilon \phi_i(x/\varepsilon) \partial_i \bar{\rho}(x) + \dots$$

and therefore that

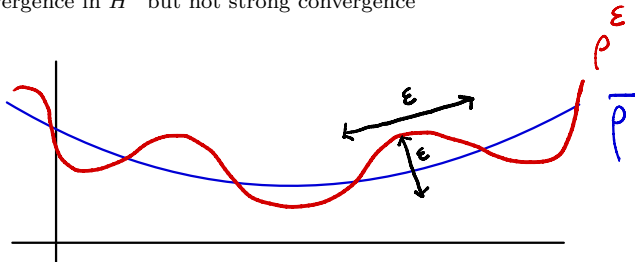
$$\nabla \rho^\varepsilon \simeq \nabla \bar{\rho} + \nabla \phi_i(x/\varepsilon) \partial_i \bar{\rho}.$$

As $\varepsilon \rightarrow 0$,

$$\nabla \rho^\varepsilon \rightharpoonup \nabla \bar{\rho} \text{ weakly in } L^2(\mathbb{R}^d \times [0, T]; \mathbb{R}^d),$$

but not strongly.

- weak convergence in H^1 but not strong convergence



IV. Periodic homogenization of divergence form equations

The *homogenization error*: the error in the two-scale expansion,

$$w^\varepsilon = \rho^\varepsilon - \bar{\rho} - \varepsilon \phi_i(x/\varepsilon) \partial_i \bar{\rho},$$

for the correctors ϕ_i and the homogenized solution $\bar{\rho}$ satisfying

$$-\nabla \cdot \bar{a} \nabla \bar{\rho} = f \text{ in } U \text{ with } \bar{\rho} = g \text{ on } \partial U \text{ and } -\nabla \cdot a(e_i + \nabla \phi_i) = 0 \text{ in } \mathbb{T}^d.$$

The homogenization error w^ε satisfies

$$\begin{aligned} -\nabla \cdot a(x/\varepsilon) \nabla w^\varepsilon &= f - \nabla \cdot (a(x/\varepsilon) e_i \partial_i \bar{\rho}) - \nabla \cdot (a(x/\varepsilon) \nabla \phi_i(x/\varepsilon) \partial_i \bar{\rho}) \\ &\quad - \varepsilon \nabla \cdot (a(x/\varepsilon) \phi_i(x/\varepsilon) \nabla (\partial_i \bar{\rho})). \end{aligned}$$

After adding and subtracting

$$-\nabla \cdot \bar{a} \nabla \bar{\rho} = f,$$

we have that

$$\begin{aligned} -\nabla \cdot a(x/\varepsilon) \nabla w^\varepsilon &= \nabla \cdot ((a(x/\varepsilon)(e_i + \nabla \phi_i(x/\varepsilon)) - \bar{a} e_i) \partial_i \bar{\rho}) \\ &\quad + \varepsilon \nabla \cdot (a(x/\varepsilon) \phi_i(x/\varepsilon) (\nabla \partial_i \bar{\rho})). \end{aligned}$$

The energy estimate, for $c = c(d, f, g, \lambda, \Lambda) \in (0, \infty)$,

$$\int_U |\nabla w^\varepsilon|^2 \leq c \left(\|(a(x/\varepsilon)(e_i + \nabla \phi_i(x/\varepsilon)) - \bar{a} e_i)\|_{L^2(\mathbb{T}^d)}^2 + \varepsilon \|\phi_i\|_{L^2(\mathbb{T}^d)}^2 \right).$$

IV. Periodic homogenization of divergence form equations

The homogenization error $w^\varepsilon = \rho^\varepsilon - \bar{\rho} - \varepsilon \phi_i(x/\varepsilon) \partial_i \bar{\rho}$ satisfies

$$-\nabla \cdot a(x/\varepsilon) \nabla w^\varepsilon = \nabla \cdot ((a(x/\varepsilon)(e_i + \nabla \phi_i(x/\varepsilon)) - \bar{a}e_i) \partial_i \bar{\rho}) \\ + \varepsilon \nabla \cdot (a(x/\varepsilon) \phi_i(x/\varepsilon) (\nabla \partial_i \bar{\rho})).$$

The essential observation is that the vectors

$$q_i = a(e_i + \nabla \phi_i) - \bar{a}e_i,$$

are mean zero:

$$\langle q_i \rangle = \langle a(e_i + \nabla \phi_i) \rangle - \bar{a}e_i = 0$$

and divergence free:

$$\nabla \cdot q_i = \nabla \cdot a(y)(e_i + \nabla \phi_i(y)) = 0.$$

- De Rham cohomology—every mean zero, divergence free field has a potential field
 - there exists σ_i such that $\nabla \cdot \sigma_i = q_i$.
- σ is a $d \times d$ skew-symmetric matrix
 - vector field is a $(d-1)$ -form, skew symmetric matrix is a $(d-2)$ -form
 - the “stream matrix” from fluids

IV. Periodic homogenization of divergence form equations

Let $q = (q_i): \mathbb{T}^d \rightarrow \mathbb{R}^d$ be mean zero and divergence free.

There exists a skew-symmetric matrix $\sigma: \mathbb{T}^d \rightarrow \mathbb{R}^{d \times d}$ such that

$$\nabla \cdot \sigma = q \text{ where } (\nabla \cdot \sigma)_i = \partial_k \sigma_{ik}. \quad (3)$$

The solution to (3) is not unique—shift by any matrix with rows that are divergence free.

A variational problem: for $\sigma = (\sigma_{jk})$ minimize the energy

$$\int_{\mathbb{T}^d} |\nabla \sigma_{jk}|^2 \text{ subject to the constraint } \nabla \cdot \sigma = q.$$

This leads to the well-posed equation

$$-\Delta \sigma_{jk} = \partial_j q_k - \partial_k q_j \text{ on } \mathbb{T}^d \text{ with } \langle \sigma_{jk} \rangle = 0.$$

We define $\sigma = (\sigma_{jk})$ and observe that

$$\Delta(\nabla \cdot \sigma)_i = \Delta(\partial_k \sigma_{ik}) = \partial_k(\Delta \sigma_{ik}) = \Delta q_i - \partial_i(\nabla \cdot q) = \Delta q_i.$$

We have that

$$\Delta((\nabla \cdot \sigma)_i - q_i) = 0 \text{ and so } (\nabla \cdot \sigma)_i - q_i = \langle (\nabla \cdot \sigma)_i - q_i \rangle = 0,$$

and, therefore,

$$\nabla \cdot \sigma = q.$$

IV. Periodic homogenization of divergence form equations

The *homogenization error*: the error in the two-scale expansion,

$$w^\varepsilon = \rho^\varepsilon - \bar{\rho} - \varepsilon \phi_i(x/\varepsilon) \partial_i \bar{\rho},$$

for the correctors ϕ_i and the homogenized solution $\bar{\rho}$ satisfying

$$-\nabla \cdot \bar{a} \nabla \bar{\rho} = f \text{ in } U \text{ with } \bar{\rho} = g \text{ on } \partial U \text{ and } -\nabla \cdot a(e_i + \nabla \phi_i) = 0 \text{ in } \mathbb{T}^d.$$

The homogenization error w^ε satisfies

$$\begin{aligned} -\nabla \cdot a(x/\varepsilon) \nabla w^\varepsilon &= \nabla \cdot ((a(x/\varepsilon)(e_i + \nabla \phi_i(x/\varepsilon)) - \bar{a}e_i) \partial_i \bar{\rho}) \\ &\quad + \varepsilon \nabla \cdot (a(x/\varepsilon) \phi_i(x/\varepsilon) (\nabla \partial_i \bar{\rho})). \end{aligned}$$

The *fluxes* q_i : divergence free fields

$$q_i = a(e_i + \nabla \phi_i) - \bar{a}e_i.$$

The *flux correctors*: skew-symmetric matrices σ_i satisfying

$$\nabla \cdot \sigma_i = q_i \text{ fixed by } -\Delta \sigma_{ijk} = \partial_j q_{ik} - \partial_k q_{ij}.$$

It follows from the skew-symmetry that distributionally

$$\nabla \cdot (q_i(x/\varepsilon) \partial_i \bar{\rho}) = -\varepsilon \nabla \cdot (\sigma_i(x/\varepsilon) \nabla \partial_i \bar{\rho}),$$

and, therefore,

$$-\nabla \cdot a(x/\varepsilon) \nabla w^\varepsilon = \varepsilon \nabla \cdot ((a \phi_i(x/\varepsilon) - \sigma_i(x/\varepsilon)) \nabla \partial_i \bar{\rho}).$$

IV. Periodic homogenization of divergence form equations

Periodic homogenization of divergence form equations

Let $a: \mathbb{T}^d \rightarrow \mathbb{R}^{d \times d}$ be bounded, measurable, and uniformly elliptic. Let $\rho_0 \in C_c^\infty(\mathbb{R}^d)$ and for every $\varepsilon \in (0, 1)$ let

$$-\nabla \cdot a(x/\varepsilon) \nabla \rho^\varepsilon = f \text{ in } U \text{ with } \rho^\varepsilon = g \text{ on } \partial U,$$

and let $\bar{\rho}$ solve

$$-\nabla \cdot \bar{a} \nabla \bar{\rho} = f \text{ in } U \text{ with } \bar{\rho} = g \text{ on } \partial U,$$

for $\bar{a}e_i = \langle a(e_i + \nabla \phi_i) \rangle$ defined by the correctors $-\nabla \cdot a(e_i + \nabla \phi_i) = 0$. Then there exists $c = c(\rho_0, d, f, \lambda, \Lambda) \in (0, \infty)$ such that the homogenization error

$$w^\varepsilon = \rho^\varepsilon - \bar{\rho} - \varepsilon \phi_i(x/\varepsilon) \partial_i \bar{\rho}$$

satisfies

$$\|w^\varepsilon\|_{H^1(U)} \leq c\varepsilon \left(\|\phi_i\|_{L^2(\mathbb{T}^d)} + \|\sigma_i\|_{L^2(\mathbb{T}^d)} \right).$$

- strong H^1 -convergence of the two-scale expansion
 - the two-scale expansion corrects the function *and* the gradient
- the a -harmonic coordinates $x_i + \phi_i(x)$ —sub-quadratic a -harmonic functions
 - Liouville theorems—these are the linear functions in the geometry of a

V. Periodic homogenization of non-divergence form equations

The diffusion equation

$$\operatorname{tr}(a(x/\varepsilon)\nabla^2\rho^\varepsilon) + \varepsilon^{-1}b(x/\varepsilon) \cdot \nabla\rho^\varepsilon = f \text{ in } U \text{ with } \rho^\varepsilon = g \text{ on } \partial U.$$

- Postulate an asymptotic expansion of the form

$$\tilde{\rho}^\varepsilon = \bar{\rho} + \varepsilon\phi_i(x/\varepsilon)\partial_i\bar{\rho} + \varepsilon^2\psi(x/\varepsilon) + \dots$$

- The *corrector equation* at order ε^{-1} :

$$\operatorname{tr}(a(y)\nabla^2\phi_i) + b(y) \cdot \nabla\phi_i = -b_i \text{ in } \mathbb{T}^d.$$

Solvability requires $\langle b, m \rangle = 0$ for the invariant measure $\mathcal{L}^*m = 0$.

- At order ε^0 the solvability condition for ψ requires that $\bar{\rho}$ solves

$$\partial_t\bar{\rho} = \operatorname{tr}(\bar{a}\nabla^2\bar{\rho}) \text{ for } \bar{a} = \int_{\mathbb{T}^d} (a(y)(1 + \nabla\phi(y)) + \phi(y) \otimes b(y)) m(y) dy.$$

- For the process $dX_t = \sigma(X_t)dB_t + b(X_t)dt$ the process

$$M_t = X_t + \phi(X_t) \text{ for } \phi = (\phi_1, \dots, \phi_d) \text{ is a martingale.}$$

- Form the decomposition

$$\varepsilon X_{t/\varepsilon^2} = \varepsilon M_{t/\varepsilon^2} - \varepsilon\phi(X_{t/\varepsilon^2}).$$

V. Periodic homogenization of non-divergence form equations

An alternate approach using the invariant measure: for $\mathcal{L}^*m = 0$,

$$\begin{aligned} & \operatorname{tr}(m(x/\varepsilon)a(x/\varepsilon)\nabla^2\rho^\varepsilon) + \varepsilon^{-1}m(x/\varepsilon)b(x/\varepsilon) \cdot \nabla\rho^\varepsilon \\ &= \nabla \cdot (m(x/\varepsilon)a(x/\varepsilon)\nabla\rho^\varepsilon) - \varepsilon^{-1}(\nabla \cdot (m(x/\varepsilon)a(x/\varepsilon)) - m(x/\varepsilon)b(x/\varepsilon)) \cdot \nabla\rho^\varepsilon. \end{aligned}$$

- *Divergence-free drift*: the vector

$$\tilde{b}(y) = \nabla \cdot (m(y)a(y)) - m(y)b(y) \text{ is divergence-free,}$$

since $\nabla \cdot \tilde{b} = \mathcal{L}^*m$.

- *Mean zero drift*: We have that

$$\langle \tilde{b} \rangle = \langle \nabla \cdot (ma) - mb \rangle = -\langle mb \rangle = 0,$$

if and only if b is perpendicular to m in the L^2 -sense that $\langle b, m \rangle_{L^2(\mathbb{T}^d)} = 0$.

- If the solvability condition $\langle b, m \rangle_{L^2(\mathbb{T}^d)} = 0$ is satisfied, there exists a potential $\tilde{\sigma}$ with

$$\nabla \cdot \tilde{\sigma} = \tilde{b},$$

such that

$$\operatorname{tr}(m(x/\varepsilon)a(x/\varepsilon)\nabla^2\rho^\varepsilon) + \varepsilon^{-1}m(x/\varepsilon)b(x/\varepsilon) \cdot \nabla\rho^\varepsilon = \nabla \cdot ((m(x/\varepsilon)a(x/\varepsilon) + \tilde{\sigma}(x/\varepsilon))\nabla\rho^\varepsilon).$$

- For the new “diffusion matrix”

$$\tilde{a} = am + \tilde{\sigma} \text{ we have } \nabla \cdot \tilde{a}(x/\varepsilon)\nabla\rho^\varepsilon = f(x)m(x/\varepsilon).$$

V. Periodic homogenization of non-divergence form equations

After multiplying the equation by $m(x/\varepsilon)$,

$$\begin{aligned} & \operatorname{tr}(m(x/\varepsilon)a(x/\varepsilon)\nabla^2\rho^\varepsilon) + \varepsilon^{-1}m(x/\varepsilon)b(x/\varepsilon) \cdot \nabla\rho^\varepsilon \\ &= \nabla \cdot (m(x/\varepsilon)a(x/\varepsilon)\nabla\rho^\varepsilon) - \varepsilon^{-1}(\nabla \cdot (m(x/\varepsilon)a(x/\varepsilon)) - m(x/\varepsilon)b(x/\varepsilon)) \cdot \nabla\rho^\varepsilon \\ &= \nabla \cdot ((m(x/\varepsilon)a(x/\varepsilon) + \tilde{\sigma}(x/\varepsilon))\nabla\rho^\varepsilon) = fm(x/\varepsilon), \end{aligned}$$

for $\langle b, m \rangle_{L^2(\mathbb{T}^d)} = 0$ and for the skew-symmetric matrix $\tilde{\sigma}$ satisfying

$$\nabla \cdot \tilde{\sigma} = \nabla \cdot (m(x/\varepsilon)a(x/\varepsilon)) - m(x/\varepsilon)b(x/\varepsilon).$$

The correctors

$$-\nabla \cdot (am + \tilde{\sigma})(e_i + \nabla\phi_i) = 0 \quad \text{in } \mathbb{T}^d.$$

Observe that by Hölder's inequality and Young's inequality that, for $c = c(d) \in (0, \infty)$,

$$\int_{\mathbb{T}^d} \langle (am + \tilde{\sigma})\nabla\phi, \nabla\phi \rangle = \int_{\mathbb{T}^d} \langle a\nabla\phi, \nabla\phi \rangle m \leq c \left(\|am\|_{L^2(\mathbb{T}^d)}^2 + \|\tilde{\sigma}\|_{L^2(\mathbb{T}^d)}^2 \right).$$

The homogenized matrix \bar{a} is

$$\bar{a}e_i := \langle (am + \tilde{\sigma})(e_i + \nabla\phi_i) \rangle,$$

and $\bar{\rho}$ solves

$$\nabla \cdot \bar{a}\nabla\bar{\rho} = fm(x/\varepsilon) \quad \text{in } U \quad \text{with } \bar{\rho} = g \quad \text{on } \partial U.$$

V. Periodic homogenization of non-divergence form equations

Periodic homogenization of non-divergence form equations

Assume that $a \in C^{1,\alpha}(\mathbb{T}^d; \mathbb{R}^{d \times d})$ is uniformly elliptic, assume that $b \in C^\alpha(\mathbb{T}^d; \mathbb{R}^d)$, and assume that $\rho_0 \in C_c^\infty(\mathbb{T}^d)$. Let m be the invariant measure $\mathcal{L}^* m = 0$ and assume that $\langle b, m \rangle_{L^2(\mathbb{T}^d)} = 0$. Let $\tilde{\sigma}$ be the vector potential satisfying

$$\nabla \cdot \tilde{\sigma} = \nabla \cdot (am) - bm,$$

and let $\tilde{a} \in C^\alpha(\mathbb{R}^d)$ be defined by $\tilde{a} = am + \tilde{\sigma}$. Define the homogenization correctors

$$-\nabla \cdot \tilde{a}(e_i + \nabla \phi_i) = 0 \text{ in } \mathbb{T}^d.$$

Then, for the effective matrix $\bar{a}e_i = \langle \tilde{a}(e_i + \nabla \phi_i) \rangle$, for $\bar{\rho}$ satisfying

$$\nabla \cdot \bar{a} \nabla \bar{\rho} = fm(x/\varepsilon) \text{ in } U \text{ with } \bar{\rho} = g \text{ on } \partial U,$$

the homogenization error

$$w^\varepsilon = \rho^\varepsilon - \bar{\rho} - \varepsilon \phi_i \partial_i \bar{\rho},$$

satisfies, for some $c = c(f, g, d, \lambda, \Lambda) \in (0, \infty)$,

$$\|w^\varepsilon\|_{L^2[0,T]; H^1(\mathbb{T}^d)} \leq c\varepsilon \left(\|\phi_i\|_{L^2(\mathbb{T}^d)} + \|\sigma_i\|_{L^2(\mathbb{T}^d)} \right),$$

for the flux correctors $\nabla \cdot \sigma_i = \tilde{a}(e_i + \nabla \phi_i) - \bar{a}e_i$.

- Energy estimates for the correctors: there exists $c = c(a, m, \tilde{\sigma}) \in (0, \infty)$ such that

$$\int_{\mathbb{T}^d} |\nabla \phi_i(y)|^2 m(y) dy \leq c.$$

V. Periodic homogenization of non-divergence form equations

Important examples with $\bar{b} = 0$:

- *Divergence-form:*

$$-\nabla \cdot a(x/\varepsilon) \nabla \rho^\varepsilon = \operatorname{tr}\left(a\left(\frac{x}{\varepsilon}\right) \nabla^2 \rho^\varepsilon\right) + \varepsilon^{-1} (\nabla \cdot a^t(x/\varepsilon)) \cdot \nabla \rho^\varepsilon.$$

In this case $m = 1$ and, as the integral of a periodic gradient,

$$\langle b, m \rangle_{L^2(\mathbb{T}^d)} = \int_{\mathbb{T}^d} (\nabla \cdot a^t(y)) \, dy = 0.$$

- *Mean-zero divergence free drift:* for a potential s with $\nabla \cdot s = b$,

$$-\nabla \cdot a(x/\varepsilon) \nabla \rho^\varepsilon + \varepsilon^{-1} b(x/\varepsilon) \cdot \nabla \rho^\varepsilon = \nabla \cdot (a + s)(x/\varepsilon) \nabla \rho^\varepsilon.$$

In this case $m = 1$ and

$$\langle b, m \rangle_{L^2(\mathbb{T}^d)} = \int_{\mathbb{T}^d} b \, dy = \int_{\mathbb{T}^d} \nabla \cdot s(y) \, dy = 0.$$

- *Brownian motion in a periodic potential:* we consider

$$-\Delta \rho^\varepsilon + \varepsilon^{-1} \nabla U(x/\varepsilon) \cdot \nabla \rho^\varepsilon = -e^{U(x/\varepsilon)} \nabla \cdot (e^{-U(x/\varepsilon)} \nabla \rho^\varepsilon).$$

The invariant measure is the Gibbs measure $m = \langle e^{-U} \rangle^{-1} e^{-U}$ and

$$\langle b, m \rangle_{L^2(\mathbb{T}^d)} = \langle e^{-U} \rangle^{-1} \int_{\mathbb{T}^d} e^{-U} \nabla U \, dy = -\langle e^{-U} \rangle^{-1} \int_{\mathbb{T}^d} \nabla e^{-U} \, dy = 0.$$

- *Symmetry:* restricted isotropy in law implies $\bar{b} = 0$.

V. Periodic homogenization of non-divergence form equations

The case $\bar{b} = \langle b, m \rangle_{L^2(\mathbb{T}^d)} \neq 0$.

- Constant coefficients:

$$dX_t = dB_t + \bar{b} dt.$$

In this case, for the diffusion beginning at zero,

$$X_t^\varepsilon = W_t^\varepsilon + \varepsilon^{-1} t \bar{b}.$$

Diffusive behavior after subtracting the “effective drift” \bar{b} :

$$X_t^\varepsilon - \varepsilon^{-1} t \bar{b} = W_t^\varepsilon.$$

- True in general: if $\bar{b} \neq 0$ then as $\varepsilon \rightarrow 0$ the process

$$X_t^\varepsilon \text{ is ballistic in the direction } \bar{b}.$$

However, after centering about this singular trajectory, as $\varepsilon \rightarrow 0$,

$$X_t^\varepsilon - \varepsilon^{-1} \bar{b} t \rightarrow \bar{\sigma} B_t \text{ in law.}$$

- Repeating the same proof:

$$\partial_t \rho^\varepsilon - \nabla \cdot \tilde{a}(x/\varepsilon) \nabla \rho^\varepsilon + \varepsilon^{-1} \bar{b} \cdot \nabla \rho^\varepsilon = f(x) m(x/\varepsilon),$$

and $\bar{\rho}^\varepsilon$ solves

$$\partial_t \bar{\rho}^\varepsilon - \nabla \cdot \bar{a} \nabla \bar{\rho}^\varepsilon + \varepsilon^{-1} \bar{b} \cdot \nabla \bar{\rho}^\varepsilon = f(x) m(x/\varepsilon).$$

Compare $\rho^\varepsilon(x - \varepsilon^{-1} \bar{b} t, t)$ to $\bar{\rho}^\varepsilon(x - \varepsilon^{-1} \bar{b} t, t)$.

VI. A random environment

The Poisson point process on \mathbb{R}^d :

- The probability space is the space of locally finite point measures

$$\Omega = \left\{ \omega = \sum_{i \in I} \delta_{x_i} : x_i \text{ are locally finite in } \mathbb{R}^d \right\},$$

with the sigma algebra \mathcal{F} generated by all maps of the form

$$\omega \rightarrow \omega(B) = \#\{i \in I : x_i \in B\} \text{ for Borel subsets } B \subseteq \mathbb{R}^d.$$

- For $\lambda \in (0, \infty)$ there exists a unique probability measure \mathbb{P}_λ on Ω satisfying:
 - For every Borel subset $B \subseteq \mathbb{R}^d$,

$$\mathbb{E}_\lambda[\omega(B)] = \lambda |B|.$$

- For every collection of bounded, disjoint subsets $B_1, \dots, B_N \subseteq \mathbb{R}^d$,
the random variables $\omega \rightarrow \omega(B_k)$ are independent.
- For every $y \in \mathbb{R}^d$ and measurable set $A \in \mathcal{F}$,

$$\mathbb{P}_\lambda(A) = \mathbb{P}_\lambda(A + y) \text{ for } A + y = \{\omega(\cdot + y) : \omega \in A\}.$$

VI. A random environment

Fix a Poisson point process $(\Omega, \mathcal{F}, \mathbb{P}_\lambda)$. We define the random coefficient field $a(x, \omega): \mathbb{R}^d \times \Omega \rightarrow \mathbb{R}^{d \times d}$, for $\omega = \sum_{i \in I} \delta_{x_i}$,

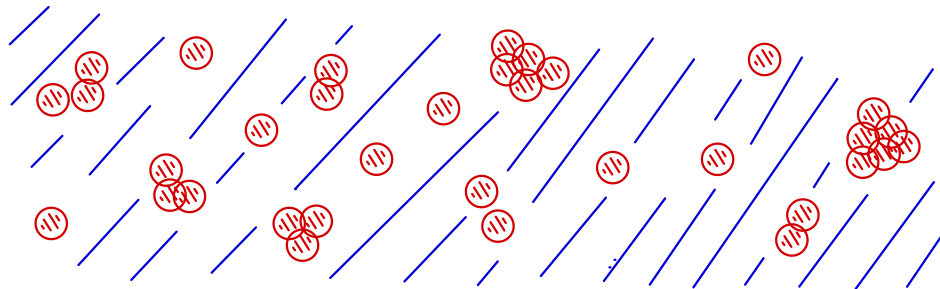
$$a(x, \omega) = \lambda_1 \mathbf{1}_{\{\cup_{i \in I} B_1(x_i)\}} + \lambda_2 \mathbf{1}_{\{\cup_{i \in I} B_1(x_i)\}^c}$$

For the measure-preserving transformation group $\{\tau_x: \mathbb{R}^d \rightarrow \mathbb{R}^d\}_{x \in \mathbb{R}^d}$ defined by

$$\tau_x(\omega)(\cdot) = \omega(\cdot - x),$$

we have

$$a(x + y, \omega) = a(y, \tau_x \omega) \text{ for every } x, y \in \mathbb{R}^d \text{ and } \omega \in \Omega.$$



VII. Stochastic homogenization

A random uniformly elliptic coefficient field $a(x, \omega): \mathbb{R}^d \times \Omega \rightarrow \mathbb{R}^{d \times d}$.

- *stationary*: for a measure-preserving transformation group $\{\tau_x: \Omega \rightarrow \Omega\}_{x \in \mathbb{R}^d}$,

$$a(x + y, \omega) = a(x, \tau_y \omega).$$

- *ergodicity*: the transformation group is qualitatively mixing, for $g: \Omega \rightarrow \mathbb{R}$,

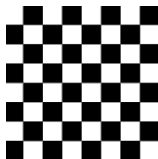
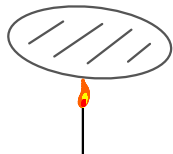
$$g(\tau_x \cdot) = g(\cdot) \text{ for every } x \in \mathbb{R}^d \text{ if and only if } g \text{ is constant.}$$

We are interested in the limiting behavior, as $\varepsilon \rightarrow 0$, of

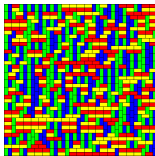
$$-\nabla \cdot a(x/\varepsilon, \omega) \nabla \rho^\varepsilon = f \text{ in } U \text{ with } \rho^\varepsilon = g \text{ on } \partial U,$$

describes a system in equilibrium:

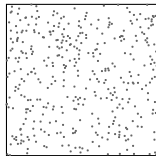
$$\oint_{B_r(x)} a(y/\varepsilon, \omega) \nabla \rho^\varepsilon(y) \cdot \nu = \int_{B_r(x)} f(y).$$



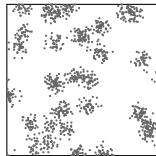
Periodic



Random tile



Poisson cloud



Cluster

VII. Stochastic homogenization

The ergodic theorem [Becker]

Let $(\Omega, \mathcal{F}, \mathbb{P})$ be a probability space equipped with an ergodic measure preserving transformation group $\{\tau_x : \Omega \rightarrow \Omega\}_{x \in \mathbb{R}^d}$.

Then for every $f \in L^1(\Omega)$, for almost every $\omega \in \Omega$,

$$\lim_{R \rightarrow \infty} \int_{B_R} f(\tau_x \omega) dx = \mathbb{E}[f].$$

And in the weak form, for every $\psi \in C_c^\infty(\mathbb{R}^d)$, as $\varepsilon \rightarrow 0$,

$$\lim_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^d} \psi(x) f(\tau_{x/\varepsilon} \omega) dx \rightarrow \mathbb{E}[f] \int_{\mathbb{R}^d} \psi(x) dx.$$

- A function $f : \mathbb{R}^d \times \Omega \rightarrow \mathbb{R}$ is *stationary* and *ergodic* if

$$f(x, \omega) = f(0, \tau_x \omega) =: g(\tau_x \omega) \text{ for every } x \in \mathbb{R}^d \text{ and } \omega \in \Omega,$$

for some measurable $g : \Omega \rightarrow \mathbb{R}$ and $\{\tau_x\}_{x \in \mathbb{R}^d}$ is ergodic.

- *Ergodicity*: large-scale spatial averages almost surely approximate the expectation:

$$\int_{B_R} f(x, \omega) dx \simeq \mathbb{E}[f] \text{ as } R \rightarrow \infty.$$

- In the weak form, almost surely,

$$f(x/\varepsilon, \omega) \rightharpoonup \mathbb{E}[f] \text{ weakly as } \varepsilon \rightarrow 0.$$

VII. Stochastic homogenization

Stochastic homogenization: for the solutions

$$-\nabla \cdot a(x/\varepsilon, \omega) \rho^\varepsilon(x, \omega) = f \text{ in } U \text{ with } \rho^\varepsilon(\cdot, \omega) = g \text{ on } \partial U,$$

there exists a deterministic $\bar{a} \in \mathbb{R}^{d \times d}$ such that

$$\rho^\varepsilon(\cdot, \omega) \rightarrow \bar{\rho} \text{ almost surely as } \varepsilon \rightarrow 0,$$

for the solution $\bar{\rho}$ of

$$-\nabla \cdot \bar{a} \nabla \bar{\rho} = f \text{ in } U \text{ with } \bar{\rho} = g \text{ on } \partial U.$$

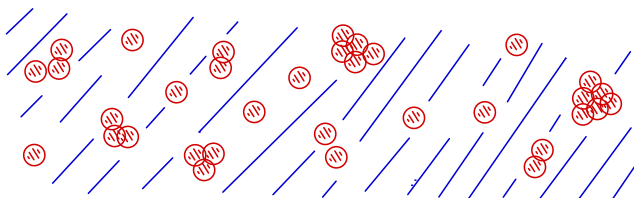
Diffusion in random environment: in the symmetric case, for the diffusion processes

$$dX_t^\omega = \sigma(X_t^\omega, \omega) dB_t + b(X_t^\omega, \omega) dt,$$

for $a = \frac{1}{2} \sigma \sigma^t$, we have almost surely that

$$\varepsilon X_{t/\varepsilon^2}^\omega \rightarrow \bar{\sigma} B_t \text{ in law,}$$

for $\bar{a} = \frac{1}{2} \bar{\sigma} \bar{\sigma}^t$.



VII. Stochastic homogenization

—The equation

$$-\nabla \cdot a(x/\varepsilon, \omega) \nabla \rho^\varepsilon = f \text{ in } U \text{ with } \rho^\varepsilon = g \text{ on } \partial U.$$

—The asymptotic expansion

$$\rho^\varepsilon(x, \omega) = \bar{\rho}(x) + \varepsilon \phi_i(x/\varepsilon, \omega) \partial_i \bar{\rho}(x) + \dots$$

—Almost surely by the ergodic theorem

$$a(x/\varepsilon, \omega)(e_i + \nabla \phi_i(x/\varepsilon, \omega)) \rightarrow \langle a(0, \omega)(e_i + \nabla \phi_i(0, \omega)) \rangle =: \bar{a} e_i.$$

—By stationarity we have that $a(x, \omega) = A(\tau_x \omega)$ and $\nabla \phi(x, \omega) = \Phi_i(\tau_x \omega)$, so that

$$a(x/\varepsilon, \omega)(e_i + \nabla \phi_i(x/\varepsilon, \omega)) \rightarrow \langle a(0, \omega)(e_i + \nabla \phi_i(0, \omega)) \rangle = \mathbb{E}[A(e_i + \Phi_i)].$$

—The first-order correctors ϕ_i almost surely satisfy on the whole space

$$-\nabla \cdot a(y, \omega)(e_i + \nabla \phi_i(y, \omega)) = 0 \text{ in } \mathbb{R}^d.$$

—For the homogenized coefficient \bar{a} we have

$$-\nabla \cdot \bar{a} \nabla \bar{\rho} = f \text{ in } U \text{ with } \bar{\rho} = g \text{ on } \partial U.$$

VII. Stochastic homogenization

The corrector equation

$$-\nabla \cdot a(y, \omega)(e_i + \nabla \phi_i(y, \omega)) = 0 \text{ in } \mathbb{R}^d.$$

The periodic case:

- The probability space is the torus,

$\Omega = \mathbb{T}^d$ with the Lebesgue sigma algebra and the normalized Lebesgue measure.

- The “random” variable

$$A: \mathbb{T}^d \rightarrow \mathbb{R}^{d \times d} \text{ is 1-periodic.}$$

- The transformation group $\{\tau_x\}_{x \in \mathbb{R}^d}$ is defined by

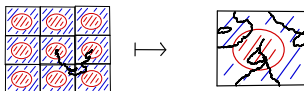
$$\tau_x \omega = x + \omega \in \mathbb{T}^d \text{ for every } x \in \mathbb{R}^d \text{ and } \omega \in \mathbb{T}^d.$$

- The stationary “random” coefficient field is

$$a(x, \omega) = A(x + \omega) = A(\tau_x \omega) \text{ for every } x \in \mathbb{R}^d \text{ and } \omega \in \mathbb{T}^d.$$

Lift the corrector equation to \mathbb{T}^d using *the environment from the point of view of the particle*:

$$-\nabla \cdot A(y)(e_i + \nabla \phi_i(y)) = 0 \text{ in } \mathbb{T}^d.$$



VII. Stochastic homogenization

The corrector equation

$$-\nabla \cdot a(y, \omega)(e_i + \nabla \phi_i(y, \omega)) = 0 \text{ in } \mathbb{R}^d.$$

and the asymptotic expansion $\rho^\varepsilon = \bar{\rho} + \varepsilon \phi_i(x/\varepsilon, \omega) \partial_i \bar{\rho} + \dots$

- Validity of the asymptotic expansion requires *sublinearity*: almost surely,

$$\lim_{\varepsilon \rightarrow 0} \left(\sup_{B_1} (\varepsilon |\phi_i(x/\varepsilon, \omega)|) \right) = \lim_{R \rightarrow \infty} \left(R^{-1} \left(\sup_{B_R} |\phi_i(y, \omega)| \right) \right) = 0.$$

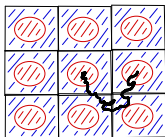
- In an L^2 -sense, almost surely,

$$\lim_{R \rightarrow \infty} \left(R^{-1} \left(\int_{B_R} \phi_i^2(y, \omega) dy \right)^{\frac{1}{2}} \right) = 0.$$

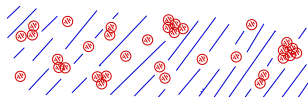
- A true *correction* of the diffusion process $dX_t^\omega = \sigma(X_t^\omega, \omega) dB_t + b(X_t^\omega, \omega) dt$:

$$X_t^\omega = X_t^\omega + \phi(X_t^\omega) - \phi(X_t^\omega).$$

- *The environment from the point of view of the particle:*



$$\bar{\omega}_t = \tau_{X_t^\omega} \omega.$$



VII. Stochastic homogenization

How to lift the corrector equation to Ω :

$$-\nabla \cdot a(y, \omega)(e_i + \nabla \phi_i(y, \omega)) = 0 \text{ in } \mathbb{R}^d.$$

- Differential operators on Ω :

$$D_i f(\omega) = \lim_{h \rightarrow 0} \frac{f(\tau_{he_i} \omega) - f(\omega)}{h} \text{ strongly in } L^2(\Omega).$$

- Smooth functions on Ω : for $\psi \in C_c^\infty(\mathbb{R}^d)$ and $f \in L^\infty(\Omega)$,

$$f_\psi(\omega) = \int_{\mathbb{R}^d} f(\tau_x \omega) \psi(x) dx.$$

- Formally we have the H^1 -space

$$\mathcal{H}^1(\Omega) = \cap_{i=1}^d \mathcal{D}(D_i).$$

- We can hope to solve

$$-D \cdot a(e_i + D\phi_i) = 0 \text{ in } \Omega,$$

in the sense that

$$\mathbb{E}[a(e_i + D\phi_i) \cdot D\psi] = 0 \text{ for all } \psi \in \mathcal{H}^1(\Omega).$$

- No compactness, no Poincaré inequality, no Fredholm alternative.

VII. Stochastic homogenization

How to lift the corrector equation to Ω :

$$-\nabla \cdot a(y, \omega)(e_i + \nabla \phi_i(y, \omega)) = 0 \text{ in } \mathbb{R}^d.$$

- The differential operators and H^1 -space:

$$D_i f(\omega) = \lim_{h \rightarrow 0} \frac{f(\tau_{he_i} \omega) - f(\omega)}{h} \text{ strongly in } L^2(\Omega) \text{ and } \mathcal{H}^1(\Omega) = \cap_{i=1}^d \mathcal{D}(D_i).$$

- The space of generalized gradients:

$$L^2_{\text{pot}}(\Omega) = \overline{\{D\psi : \psi \in \mathcal{H}^1(\Omega)\}}^{L^2(\Omega; \mathbb{R}^d)}.$$

- Every $\Phi \in L^2_{\text{pot}}(\Omega)$ is a gradient in the sense that it is distributionally curl free:

$$D_i \Phi_j = D_j \Phi_i \text{ for every } i, j \in \{1, \dots, d\}.$$

- The Lax-Milgram lemma: there exists a unique $\Phi_i \in L^2_{\text{pot}}(\Omega)$ satisfying

$$-D \cdot A(e_i + \Phi_i) = 0 \text{ in } \Omega,$$

in the sense that

$$\mathbb{E} [A(e_i + \Phi_i) \cdot \Psi] = 0 \text{ for every } \Psi \in L^2_{\text{pot}}(\Omega).$$

- We construct the *stationary gradient* of the corrector.

VII. Stochastic homogenization

- We construct the stationary gradients of the correctors and flux correctors:

$$-D \cdot A(e_i + \Phi_i) = 0 \quad \text{and} \quad -D \cdot \Sigma_{ijk} = D_j Q_{ik} - D_k Q_{ij},$$

for the fluxes $Q_i = A(e_i + \Phi_i) - \bar{a}e_i$.

- The correctors and flux correctors are almost surely defined by

$$\int_{B_1} \phi_i(y, \omega) dy = 0 \quad \text{with} \quad \nabla \phi_i(x, \omega) = \Phi_i(\tau_x \omega),$$

and

$$\int_{B_1} \sigma_{ijk}(y, \omega) dy = 0 \quad \text{with} \quad \nabla \sigma_{ijk}(x, \omega) = \Sigma_{ijk}(\tau_x \omega).$$

- For the fluxes $q_i = a(e_i + \nabla \phi_i) - \bar{a}e_i$ and for $\sigma_i = (\sigma_{ijk})$, almost surely,

$$-\nabla \cdot a(y, \omega)(e_i + \nabla \phi_i(y, \omega)) = 0 \quad \text{and} \quad \nabla \cdot \sigma_i(y, \omega) = q_i(y, \omega) \quad \text{on} \quad \mathbb{R}^d.$$

- Almost surely the homogenization error

$$w^\varepsilon(x, \omega) = \rho^\varepsilon(x, \omega) - \bar{\rho}(x) - \varepsilon \phi_i(x/\varepsilon, \omega) \partial_i \bar{\rho}$$

solves the equation

$$-\nabla \cdot a(x/\varepsilon, \omega) \nabla w^\varepsilon = \varepsilon \nabla \cdot ((a(x/\varepsilon, \omega) \phi_i(x/\varepsilon, \omega) - \sigma_i(x/\varepsilon, \omega)) \nabla \partial_i \bar{\rho}).$$

VII. Stochastic homogenization

- Almost surely the homogenization error

$$w^\varepsilon(x, \omega) = \rho^\varepsilon(x, \omega) - \bar{\rho}(x) - \varepsilon \phi_i(x/\varepsilon, \omega) \partial_i \bar{\rho}$$

solves the equation

$$-\nabla \cdot a(x/\varepsilon, \omega) \nabla w^\varepsilon = \varepsilon \nabla \cdot ((a(x/\varepsilon, \omega) \phi_i(x/\varepsilon, \omega) - \sigma_i(x/\varepsilon, \omega)) \nabla \partial_i \bar{\rho}).$$

- The energy estimate, for some $c = c(\lambda, \Lambda, d, f, g) \in (0, \infty)$,

$$\int_U |\nabla w^\varepsilon|^2 \leq c \left(\int_U |\varepsilon \phi_i(x/\varepsilon, \omega)|^2 + |\varepsilon \sigma_i(x/\varepsilon, \omega)|^2 \right).$$

- Homogenization requires L^2 -sublinearity:

$$\lim_{\varepsilon \rightarrow 0} \left(\int_U |\varepsilon \phi_i(x/\varepsilon, \omega)|^2 dx \right) = 0.$$

- It suffices to prove almost surely that

$$\lim_{\varepsilon \rightarrow 0} \left(\int_U |\varepsilon \phi_i(x/\varepsilon, \omega) - \langle \varepsilon \phi_i(\cdot/\varepsilon, \omega) \rangle_U|^2 dy \right) = 0,$$

for $\langle \varepsilon \phi_i(\cdot/\varepsilon, \omega) \rangle_U = \int_U \varepsilon \phi_i(x/\varepsilon, \omega) dx$.

VII. Stochastic homogenization

To prove that:

$$\lim_{\varepsilon \rightarrow 0} \left(\int_U |\varepsilon \phi_i(x/\varepsilon, \omega) - \langle \varepsilon \phi_i(\cdot/\varepsilon, \omega) \rangle_U|^2 dy \right) = 0.$$

- *The Poincaré inequality:* for every $\varepsilon \in (0, 1)$, for $c = c(U) \in (0, \infty)$,

$$\int_U |\varepsilon \phi_i(x/\varepsilon, \omega) - \langle \varepsilon \phi_i(\cdot/\varepsilon, \omega) \rangle_U|^2 dy \leq c \int_U |\nabla \phi_i(x/\varepsilon, \omega)|^2 dy.$$

- *The ergodic theorem:* almost surely, for $c = c(\lambda, \Lambda) \in (0, \infty)$,

$$\lim_{\varepsilon \rightarrow 0} \int_U |\nabla \phi_i(x/\varepsilon, \omega)|^2 dy = \mathbb{E} [|\Phi_i|^2] \leq c.$$

- *The Poincaré inequality:* almost surely,

$$\{\varepsilon \phi_i(x/\varepsilon, \omega) - \langle \varepsilon \phi_i(\cdot/\varepsilon, \omega) \rangle_U\}_{\varepsilon \in (0, 1)} \text{ is bounded in } H^1(U).$$

- *The ergodic theorem:* almost surely,

$$\nabla \phi_i(x/\varepsilon, \omega) \rightharpoonup \mathbb{E} [\Phi_i] = 0 \text{ weakly in } H^1(U),$$

and, therefore,

$$\varepsilon \phi_i(x/\varepsilon, \omega) - \langle \varepsilon \phi_i(\cdot/\varepsilon, \omega) \rangle_U \rightharpoonup c = 0 \text{ weakly in } H^1(U).$$

- *The Sobolev embedding theorem:* almost surely,

$$\varepsilon \phi_i(x/\varepsilon, \omega) - \langle \varepsilon \phi_i(\cdot/\varepsilon, \omega) \rangle_U \rightarrow 0 \text{ strongly in } L^2(U).$$

VII. Stochastic homogenization

Stochastic homogenization [Kozlov, Papanicolaou, Varadhan...]

Let $(\Omega, \mathcal{F}, \mathbb{P})$ be a probability space, let $a: \Omega \times \mathbb{R}^d \rightarrow \mathbb{R}^{d \times d}$ be a uniformly elliptic, stationary, and ergodic coefficient field and for every $\omega \in \Omega$ let ρ^ε solve the equation

$$-\nabla \cdot a(x/\varepsilon, \omega) \nabla \rho^\varepsilon = f \text{ in } U \text{ with } \rho^\varepsilon = g \text{ on } \partial U.$$

Let the homogenized coefficient $\bar{a} \in \mathbb{R}^{d \times d}$ be defined by

$$\bar{a}e_i := \mathbb{E}[A(e_i + \Phi_i)] \text{ for } \Phi_i \in L^2_{\text{pot}}(\Omega) \text{ satisfying } -D \cdot A(e_i + \Phi_i) = 0,$$

and let $\bar{\rho}$ be defined the homogenized solution

$$-\nabla \cdot \bar{a} \nabla \bar{\rho} = f \text{ in } U \text{ with } \bar{\rho} = g \text{ on } \partial U.$$

Then for the homogenization correctors ϕ_i defined by

$$\int_{B_1} \phi_i(y, \omega) dy = 1 \text{ with } \nabla \phi_i(y, \omega) = \Phi_i(\tau_y \omega),$$

the two-scale expansion

$$w^\varepsilon(x, \omega) = \rho^\varepsilon(x, \omega) - \bar{\rho}(x) - \varepsilon \phi_i(x/\varepsilon, \omega) \partial_i \bar{\rho}(x),$$

almost surely satisfies

$$\lim_{\varepsilon \rightarrow 0} \|w^\varepsilon\|_{H^1(U)} = 0.$$

- *A regularity theory for random elliptic operators*; Gloria, Neukamm, Otto
- *Quantitative Stochastic Homogenization and Large-Scale Regularity*; Armstrong, et al.

VII. Stochastic homogenization

Divergence-free environments [Avelleneda, Komoroski, Majda, Olla, Kozma, Tóth, F...]

Let $(\Omega, \mathcal{F}, \mathbb{P})$ be a probability space, let $a: \Omega \times \mathbb{R}^d \rightarrow \mathbb{R}^{d \times d}$ be a uniformly elliptic, stationary, and ergodic coefficient field and for every $\omega \in \Omega$ let ρ^ε solve the equation

$$-\nabla \cdot a(x/\varepsilon, \omega) \nabla \rho^\varepsilon + \varepsilon^{-1} b(x/\varepsilon, \omega) = f \text{ in } U \text{ with } \rho^\varepsilon = g \text{ on } \partial U,$$

for a stationary and ergodic, mean zero and divergence free drift $b(x, \omega) = B(\tau_x \omega)$. Assume that b admits a stationary L^p -integrable stream matrix S :

$$D \cdot S = B \text{ on } \Omega \text{ with } S \in L^p(\Omega; \mathbb{R}^{d \times d}).$$

Let the homogenized coefficient $\bar{a} \in \mathbb{R}^{d \times d}$ be defined by

$$\bar{a} e_i := \mathbb{E} [(A + S)(e_i + \Phi_i)] \text{ for } \Phi_i \in L^2_{\text{pot}}(\Omega) \text{ satisfying } -D \cdot (A + S)(e_i + \Phi_i) = 0,$$

and let $\bar{\rho}$ be the homogenized solution

$$-\nabla \cdot \bar{a} \nabla \bar{\rho} = f \text{ in } U \text{ with } \bar{\rho} = g \text{ on } \partial U.$$

If $p = 2$ then almost surely

$$\rho^\varepsilon \rightharpoonup \bar{\rho} \text{ weakly in } H^1(U).$$

If $p = d \wedge (2 + \delta)$ then the two-scale expansion $w^\varepsilon = \rho^\varepsilon - \bar{\rho} - \varepsilon \phi_i \partial_i \bar{\rho}$ almost surely satisfies

$$\lim_{\varepsilon \rightarrow 0} \|w^\varepsilon\|_{H^1(U)} = 0.$$

- Also the case $b = \nabla U$ for a stationary potential U : Gibbs measure $m = \mathbb{E} [e^{-U}]^{-1} e^{-U}$

VII. Stochastic homogenization

The diffusion $dX_t = \sigma(X_t, \omega) dB_t$.

Homogenization of balanced environments [Papanicolaou, Varadhan]

Assume that $a: \mathbb{R}^d \times \Omega \rightarrow \mathbb{R}^{d \times d}$ is uniformly elliptic, stationary, and ergodic. Then there exists a unique mutually absolutely continuous invariant measure π for the environment from the point of view of the particle on $(\Omega, \mathcal{F}, \mathbb{P})$: for every $f \in L^\infty(\Omega)$,

$$\mathbb{E}_\pi [E_{0, \omega} [f(\tau_{X_t} \omega)]] = \mathbb{E}_\pi [f].$$

The homogenized coefficient $\bar{a} \in \mathbb{R}^{d \times d}$ is defined by

$$\bar{a} = \mathbb{E}_\pi [a].$$

The solutions

$$\partial_t \rho^\varepsilon = \operatorname{tr}(a(x/\varepsilon, \omega)) \text{ on } \mathbb{R}^d \times (0, \infty) \text{ with } \rho^\varepsilon(\cdot, 0) = \rho_0,$$

converge almost surely as $\varepsilon \rightarrow 0$ to the solution

$$\partial_t \bar{\rho} = \operatorname{tr}(\bar{a} \nabla^2 \bar{\rho}) \text{ on } \mathbb{R}^d \times (0, \infty) \text{ with } \rho(\cdot, 0) = \rho_0.$$

- The Aleksandrov-Bakelman-Pucci estimate: suppose that ρ^ε solves

$$\operatorname{tr}(a(x/\varepsilon, \omega) \nabla^2 \rho^\varepsilon) = f \text{ in } B_1 \text{ with } \rho^\varepsilon = 0 \text{ on } \partial B_1.$$

Then, for $c = c(\lambda, \Lambda, d) \in (0, \infty)$ independent of $\varepsilon \in (0, 1)$,

$$\|\rho^\varepsilon\|_{L^\infty(B_1)} \leq c \|f\|_{L^d(B_1)}.$$

VII. Stochastic homogenization

Consider the diffusion in random environment

$$dX_t = \sigma(X_t, \omega) dB_t + b(X_t, \omega) dt.$$

In the periodic case, for the invariant measure m and in the central limit scaling,

$$\langle b, m \rangle_{L^2(\mathbb{T}^d)} = \bar{b} = 0 \text{ implies a diffusive behavior,}$$

and

$$\bar{b} \neq 0 \text{ implies ballistic behavior in direction } \bar{b}.$$

In the absence of an invariant measure try to rule out ballistic behavior using *symmetry*. Assume that, for every orthogonal transformation r that preserves the coordinate axis,

$$(r\sigma(x, \omega), rb(x, \omega))_{x \in \mathbb{R}^d} \text{ and } (\sigma(rx, \omega), b(rx, \omega))_{x \in \mathbb{R}^d} \text{ have the same law.}$$

Then, since for every orthogonal transformation r preserving the coordinate axis,

$$X_t \text{ and } rX_t \text{ have the same law under } \mathbb{P} \times P_{0, \omega}.$$

In the *annealed sense* we have that

$$\mathbb{E} [E_{0, \omega} [X_t]] = 0.$$

VII. Stochastic homogenization

Homogenization of isotropic diffusions [Sznitman, Zeitouni, F.]

Let $a: \mathbb{R}^d \times \Omega \rightarrow \mathbb{R}^{d \times d}$ and $b: \mathbb{R}^d \times \Omega \rightarrow \mathbb{R}^d$ be uniformly elliptic, bounded, Lipschitz continuous, and stationary coefficient fields satisfying a finite range of dependence. Assume that for every orthogonal transformation r preserving the coordinate axis

$$(ra(x, \omega)r^t, rb(x, \omega))_{x \in \mathbb{R}^d} \text{ and } (a(rx, \omega), b(rx, \omega))_{x \in \mathbb{R}^d} \text{ have the same law.}$$

Then there exists $\eta \in (0, \infty)$ such that if

$$|a - I| \leq \eta \text{ and } |b| \leq \eta,$$

then there exists $\bar{a} \in \mathbb{R}$ such that the solutions

$$\partial_t \rho^\varepsilon = \text{tr}(a(x/\varepsilon, \omega) \nabla^2 \rho^\varepsilon) + \varepsilon^{-1} b(x/\varepsilon, \omega) \cdot \nabla \rho^\varepsilon \text{ in } \mathbb{R}^d \times (0, \infty) \text{ with } \rho^\varepsilon(\cdot, 0) = \rho_0,$$

converge almost surely as $\varepsilon \rightarrow 0$ to the solution of

$$\partial_t \bar{\rho} = \bar{a} \Delta \bar{\rho} \text{ in } \mathbb{R}^d \times (0, \infty) \text{ with } \bar{\rho}(\cdot, 0) = \rho_0.$$

Furthermore, there exists a unique mutually absolutely continuous invariant measure π on $(\Omega, \mathcal{F}, \mathbb{P})$ for the process from the point of view of the particle: for every $f \in L^\infty(\Omega)$,

$$\mathbb{E}_\pi [E_{0, \omega} [f(\tau_{X_t} \omega)]] = \mathbb{E}_\pi [f].$$

- the perturbation says that for short times the process is like a Brownian motion
- an inductive renormalization argument controls traps / localization / coupling