Constructions of Strict Lyapunov Functions: Stability, Robustness, Delays, and State Constraints

Matrosov's Approach

Michael Malisoff

Strict and nonstrict Lyapunov functions

- Strict and nonstrict Lyapunov functions
- Input-to-state stability and point stabilization

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- Strictification to certify good performance

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M. Malisoff and F. Mazenc. Constructions of Strict Lyapunov Functions. Communications and Control Engineering Series, Springer-Verlag London Ltd., London, UK, 2009.

A Lyapunov function for a system $\dot{x} = \mathcal{F}(t,x)$ with state space \mathcal{X} is a positive definite proper function $V:[0,\infty)\times\mathcal{X}\to[0,\infty)$ such that $\dot{V}(t,x):=V_t(t,x)+V_x(t,x)\mathcal{F}(t,x)\leq 0$ on $[0,\infty)\times\mathcal{X}$.

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For example, $V(x) = \ln(1 + x^2)$ is a Lyapunov function for $\dot{x} = -x/(1 + x^2)$ because $\dot{V} \le -x^2/(1 + x^2)^2$, which gives global asymptotic stability, i.e., attractivity and local stability.

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However, for each constant $\bar{\delta} > 0$, we can find an x_0 such that the trajectory for $\dot{x} = -x/(1+x^2) + \bar{\delta}$ starting at $x(0) = x_0$ is unbounded, which means we lack input-to-state stability.

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Using LaSalle Invariance, we can often use nonstrict ones to prove GAS, e.g., for $\dot{x} = f(x)$ where $\dot{V}(x) := \nabla V(x) f(x)$.

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If *V* is a nonstrict Lyapunov function such that the only solution that remains in $\{x : \dot{V}(x) = 0\}$ is x = 0, then conclude GAS to 0.

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, $\dot{x}_2 = -x_1 - x_2^3$. Use $V(x) = 0.5|x|^2$. Then $\dot{V} = -x_2^4$.

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For example, take $\dot{x}_1 = x_2$, $\dot{x}_2 = -x_1 - x_2^3$. Use $V(x) = 0.5|x|^2$. Then $\dot{V} = -x_2^4$. The largest invariant set in $\{x : x_2 = 0\}$ is $\{0\}$.

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However, explicit strict Lyapunov function *constructions* are often needed in applications to certify robustness.

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This has led to significant research on explicitly constructing strict Lyapunov functions.

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We assume standard assumptions on the dynamics which hold under smooth forward completeness and time-periodicity.

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We say that $\dot{x}=\mathcal{F}(t,x,d)$ is ISS provided there exist functions $\beta\in\mathcal{KL}$ and $\gamma\in\mathcal{K}_{\infty}$ and $\bar{\alpha}\in\mathcal{K}_{\infty}$ s.t. for all initial conditions $x(t_0)=x_0\in\mathcal{X}$ and all disturbances d, the corresponding trajectories $t\mapsto \zeta(t;t_0,x_0,d)$ satisfy

$$|\zeta(t;t_0,x_0,d)| \leq \beta \left(\bar{\alpha}(|x_0|),t-t_0\right) + \gamma(|d|_{\infty}) \quad \forall t \geq t_0.$$
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UGAS: Special case where d = 0.

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Then

$$\dot{x} = f(t,x) + g(t,x) \left[K(t,x) - D_x V(t,x) \cdot g(t,x) + d \right]$$
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is ISS with respect to actuator errors d.

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Need K(t, x) and $D_x V(t, x) \cdot g(t, x)$.

Assume $\dot{x} = f(x)$ has a nonstrict Lyapunov function V so that:

$$\exists N_* > 0 \text{ s.t. } \forall q \in \mathbb{R}^n \setminus \{0\}, \ \exists i \in [1, N_*] \text{ s.t. } L_t^i V(q) \neq 0.$$
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$$L_f V = (\nabla V)f, L_f^i V = L_f(L_f^{i-1} V).$$

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In fact, if $L_f V(x(t,x_0)) \equiv 0$ along some trajectory, then $L_f^k V(x(t,x_0)) \equiv 0$ for all $t \geq 0$ and $k \in \mathbb{N}$, so $L_f^k V(x_0) \equiv 0$.

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Q: Can we transform V into a strict Lyapunov function?

A: Yes, and we can allow time varying systems and relax NDC.

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 (NDC) $L_f V = (\nabla V)f, \ L_f^i V = L_f(L_f^{i-1} V).$ Then GAS holds.

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A: Yes, and we can allow time varying systems and relax NDC.

Let $V \in C^{\infty}$ be a nonstrict Lyapunov function for $\dot{x} = f(t, x)$, $x \in \mathbb{R}^n$, with f and V having period T in t.

Assume $\dot{x} = f(x)$ has a nonstrict Lyapunov function V so that:

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 (NDC) $L_f V = (\nabla V)f, \ L_f^i V = L_f(L_f^{i-1} V).$ Then GAS holds.

In fact, if $L_t V(x(t,x_0)) \equiv 0$ along some trajectory, then $L_t^k V(x(t,x_0)) \equiv 0$ for all $t \geq 0$ and $k \in \mathbb{N}$, so $L_t^k V(x_0) \equiv 0$.

Q: Can we transform V into a strict Lyapunov function?

A: Yes, and we can allow time varying systems and relax NDC.

Let $V \in C^{\infty}$ be a nonstrict Lyapunov function for $\dot{x} = f(t, x)$, $x \in \mathbb{R}^n$, with f and V having period T in t. Goal:

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Theorem 1 (MM-FM, TAC'10)

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Then we can explicitly determine functions \mathcal{F}_j and \mathcal{G} such that

$$V^{\sharp}(t,x) = \sum_{j=1}^{\ell-1} \mathcal{F}_j(V(t,x)) A_j(t,x) + \mathcal{G}(t,V(t,x))$$
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is a strict Lyapunov function, giving UGAS of the dynamics.

Second Construction for $\dot{x} = f(x)$, $x \in \mathcal{X}$

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Assumption A There are functions h_j such that $h_j(0) = 0$ for all j; everywhere positive functions r_1, \ldots, r_m and ρ ; a proper positive definite function $V_1 : \mathcal{X} \to [0, \infty)$; and an integer N > 0 for which

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Also, $f \in C^{\infty}(\mathbb{R}^n)$, and V_1 has a positive definite quadratic lower bound in some neighborhood of $0 \in \mathbb{R}^n$.

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One can determine explicit functions $k_{\ell}, \Omega_{\ell} \in \mathcal{K}_{\infty} \cap C^{1}$ such that

$$S(x) = \sum_{\ell=1}^{N} \Omega_{\ell} \left(k_{\ell}(V_1(x)) + V_{\ell}(x) \right)$$
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Significance: New theorem says which functions V_i to pick.

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Significance: Readily extends to time periodic t-v systems.

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Change coordinates and rescale to get the error dynamics

$$\begin{cases} \dot{\tilde{x}} = -[\tilde{x} + \alpha \tilde{y}](\tilde{x} + x_*) \\ \dot{\tilde{y}} = \alpha \tilde{x}(\tilde{y} + y_*), \end{cases}$$
(12)

with state space $\mathcal{X} = (-x_*, +\infty) \times (-y_*, +\infty)$,

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Assume $\alpha > d$. Want a global strict Lyapunov function for (12).

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$$V_1(\tilde{x}, \tilde{y}) = \tilde{x} - x_* \ln\left(1 + \frac{\tilde{x}}{x_*}\right) + \tilde{y} - y_* \ln\left(1 + \frac{\tilde{y}}{y_*}\right)$$
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$$S(\tilde{x}, \tilde{y}) = V_2(\tilde{x}, \tilde{y}) + \int_0^{V_1(\tilde{x}, \tilde{y})} \phi_1(r) dr + \left[p_1(V_1(\tilde{x}, \tilde{y})) + 1 \right] V_1(\tilde{x}, \tilde{y}),$$
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Along the trajectories of the L-V error dynamics,

$$\dot{S}(t,x) \leq -\frac{1}{4} \left[\tilde{x}^2 + \left\{ (\tilde{x} + \alpha \tilde{y})(\tilde{x} + x_*) \right\}^2 \right]. \tag{16}$$

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Thank you for your attention and interest!