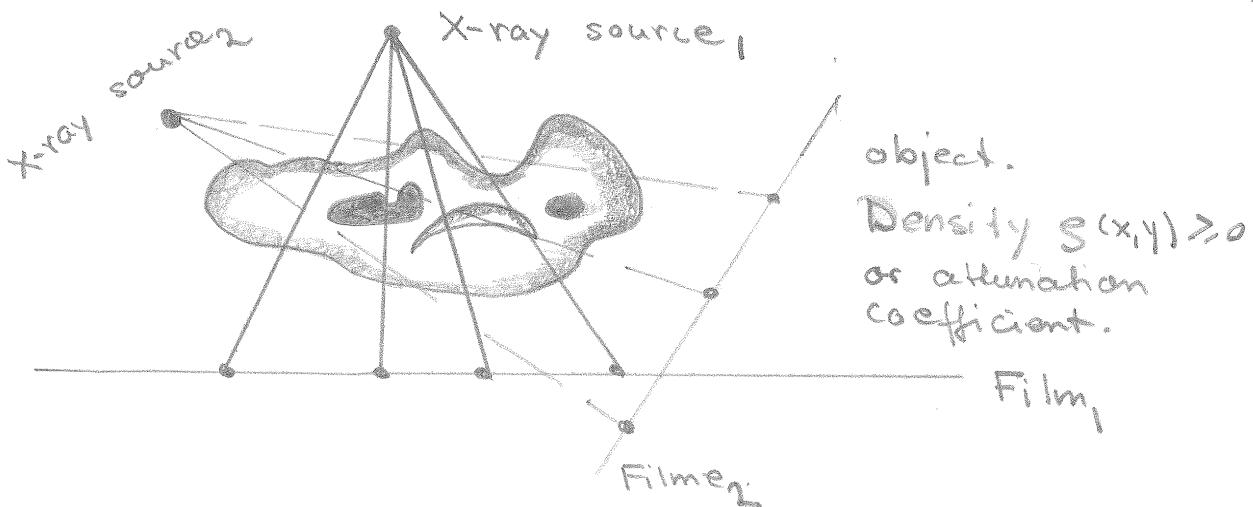


# THE RADON TRANSFORM - AND APPLICATIONS

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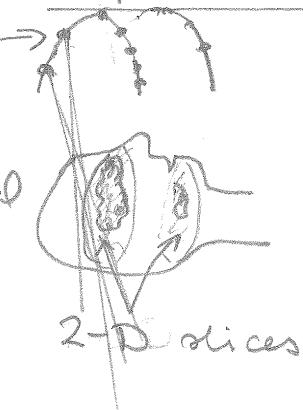
## §1 A SIMPLE MODEL

The Radon Transform has become popular because of its application in image reconstruction, in particular x-ray tomography.



The object that we are interested in is described by a function  $g : \mathbb{R}^3 \rightarrow \mathbb{R}$  (or  $\mathbb{R}^2 \rightarrow \mathbb{R}$ ), called the attenuation coefficient. It describes the absorption (or scattering) of x-rays of given energy. The first assumption is, that  $\mu = 0$  outside the body. The aim is to reconstruct  $\mu$  from the measurements. One older way to reconstruct the function  $\mu(x, y, z)$  of 3-variables by two dimensional measurement is to fix

X-ray source  
can be rotated a fixed maximal amount  
and then translated



the  $z$ -direction and consider  $f(x, y) = \mu(x, y, z)$  for several values of  $z$ .

The assumptions that are made are:

- X-rays travel along straight lines;
- All the waves are of the same frequency
- The intensity,  $I$ , of the x-ray beam satisfies Beer's law

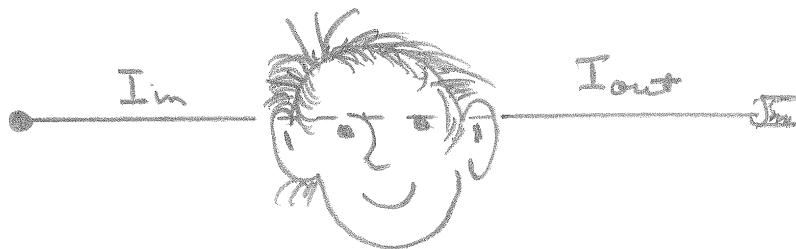
$$(x) \quad \frac{dI}{ds} = -\rho(x)I \quad L$$

where  $s$  is the arc-length along the line given by (1)

From (x) it follows, that

$$\log\left(\frac{I_{in}}{I_{out}}\right) = \int_L \rho(x)dx =: \hat{\rho}(L)$$

where  $I_{in}$  denotes the intensity of the beam before entering the body and  $I_{out}$  is the measured intensity after the beam leaves the body.



The integral transform  $\rho \mapsto \hat{\rho}$  is exactly the Radon transform introduced by Radon in 1917. It transforms "suitable" functions on  $\mathbb{R}^d$  to functions on the space of all affine lines (or more generally  $k$ -dimensional hypersurfaces in  $\mathbb{R}^d$ ).

It should be noted, that the 3 assumptions that have been made do not quite reflect the true situation, in particular the x-ray is never monochromatic, even if it might be localized around certain spectrum. Denote by  $f(E)$  the spectral function. Then a more correct model is given by the following:

$$\int_0^\infty f(E) e^{-\int_L \mu(x, E) dx} dE$$

which is non-linear in  $\mu$ !

There are other forms of applications like emission tomography

$$T_p f(L) = \int_L f(x) e^{-\int_x^L \mu(y) dy} dx$$

around. I refer to the forthcoming volume in "Proceedings of Symposia in Applied Mathematics" for further discussion.

It is also a "fact of life" that we can only use finitely many measurements to reconstruct  $\mu$  from  $\hat{\mu}$ , but  $\mu$  is never determined by finitely many values of  $\hat{\mu}$ !

This brings in topics like approximation and sampling.

If time allows, then we will also discuss 3D-versions like cone-beam tomography.

## §2 The Radon Transform

Let  $S(\mathbb{R}^n)$  be the space of rapidly decreasing functions on  $\mathbb{R}^n$ , i.e. for each  $N \in \mathbb{N}$

$$\sup_{x \in \mathbb{R}^n} (1+|x|^2)^N |f(x)| \leq C_{N,f}.$$

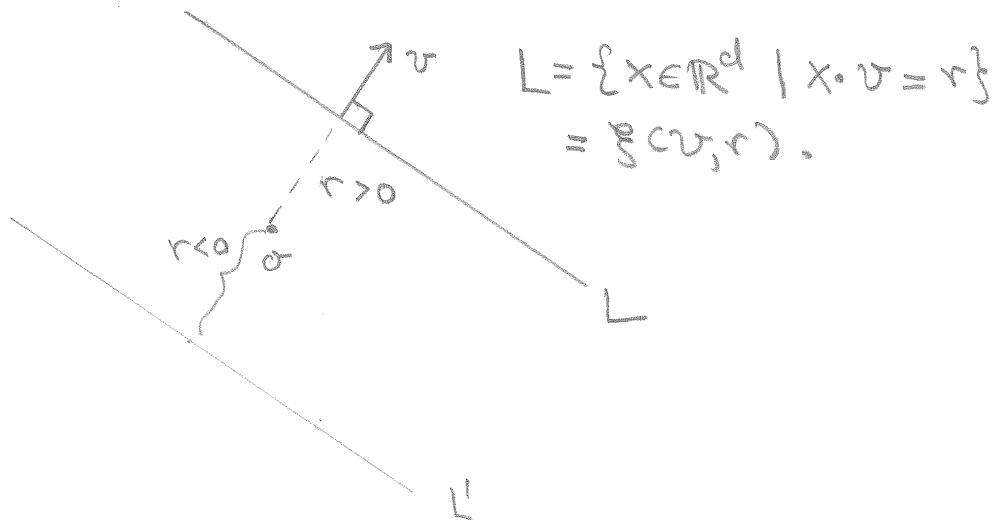
That is  $f(x) \rightarrow 0$ ,  $x \rightarrow \infty$ , faster than any  $|x|^{-N}$ . Let  $L^p(\mathbb{R}^n)$  be the space of measurable functions  $f: \mathbb{R}^n \rightarrow \mathbb{C}$  a.s.

$$\|f\|_p := \left( \int_{\mathbb{R}^n} |f(x)|^p dx \right)^{1/p} < \infty.$$

Let  $\mathcal{E}$  be the manifold of  $(d(d-1))$ -dimensional affine hyperplanes in  $\mathbb{R}^d$ . In 2-dimension this is just the set of lines in  $\mathbb{R}^d$ . A hyperplane is determined by its normal vector

$$v \in S^{d-1} = \{x \in \mathbb{R}^d \mid \|x\|=1\}$$

and its signed distance from the origin.



Note that

$$g(v, r) = g(-v, -r)$$

Thus  $\Xi = S^{d-1} \times_{\mathbb{Z}_2} \mathbb{R}$  where the subscript  $\mathbb{Z}_2$  indicates that  $(v, r)$  and  $-(v, r) = (-v, -r)$  leads to the same hyperplane. Thus we view functions on  $\Xi$  as even functions on  $S^{d-1} \times \mathbb{R}$ , i.e.  $g(v, r) = g(-v, -r)$ . The Radon transform  $f \mapsto \hat{f} = Rf$  maps functions on  $\mathbb{R}^d$  linearly to functions on  $\Xi$ :

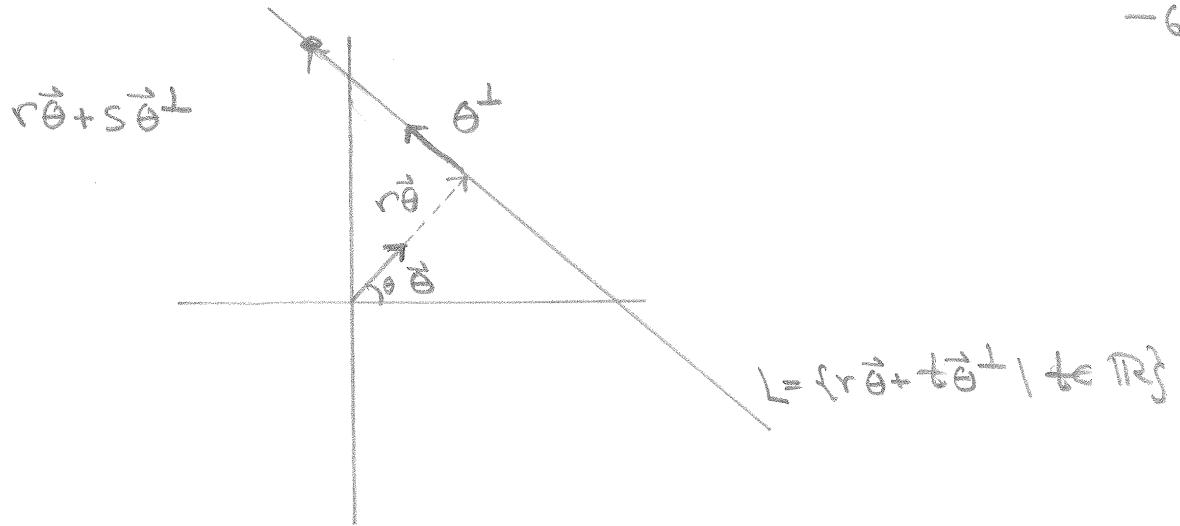
$$\begin{aligned}\hat{f}(L) &= \hat{f}(g(v, r)) = \\ &= \int_L f(x) dx \\ &= \int_{x \cdot v = r} f(x) dx, \quad f \in S(\mathbb{R}^d).\end{aligned}$$

In 2-dimensions we can write this more explicitly as:

$$(-\sin\theta, \cos\theta) = v^\perp$$
 $v = (\cos\theta, \sin\theta)$

$\vec{\theta} = v = (\cos\theta, \sin\theta)$ . Then  $v^\perp = \vec{\theta}^\perp = (-\sin\theta, \cos\theta)$  is orthogonal to  $\vec{\theta}$  and

$$L = \{r\vec{\theta} + t\vec{\theta}^\perp \mid t \in \mathbb{R}\}$$



We have therefore

$$\hat{f}(L) = \hat{f}(\vec{\theta}, r) = \int_{-\infty}^{\infty} f(r\vec{\theta} + t\vec{\theta}^\perp) dt.$$

[We have also the adjoint  
X-Rays transform  $\mathcal{P}f(\theta, x) = \int_{-\infty}^{\infty} f(x + t\vec{\theta}) dt, x \in \mathbb{S}_+^{d-1} \oplus e S^{d-1}$ .]

Let  $d\sigma(v)$  be the normalized surface measure on  $S^{d-1}$ .

In 2D this is just:

$$\begin{aligned} \int_S g(\xi) d\sigma(\xi) &= \frac{1}{2\pi} \int_0^{2\pi} g((\cos\theta, \sin\theta)) d\theta \\ &= \frac{1}{2\pi} \int_0^{2\pi} g(\vec{\theta}) d\theta. \end{aligned}$$

We can then integrate function on  $\Xi$  by

$$\begin{aligned} \int_E g(\xi) d\xi &:= \frac{1}{2} \int_{-\infty}^{\infty} \int_{S^{d-1}} g(v, r) d\sigma(v) dr \\ &= \int_0^{\infty} \int_{S^{d-1}} g(v, r) d\sigma(v) dr, \end{aligned}$$

The following questions are now natural.

- (1) Injectivity of the Radon Transform (on given function spaces)
- (2) The Image of  $\mathcal{R}$
- (3) Inversion formula.

### §3 The Central Slice Theorem and the injectivity

Define the Fourier transform of  $f \in L^1(\mathbb{R}^d)$  by

$$\mathcal{F}_d f(\lambda) = \mathcal{F}f(\lambda) = \frac{1}{(2\pi)^d} \int_{\mathbb{R}^d} f(x) e^{-i\lambda \cdot x} dx$$

$$= (f, e_\lambda)$$

$$\text{where } e_\lambda(x) = e^{-i\lambda \cdot x}.$$

(also valid for  $f \in L^1$ )

Theorem 1 (The Central Slice Theorem) Let  $f \in S(\mathbb{R}^d)$ ,  $r \in \mathbb{R}$  and  $\omega \in S^{d-1}$ . Then

$$\begin{aligned} \mathcal{F}_d f(r\omega) &= \frac{1}{(2\pi)^d} \int_{\mathbb{R}^d} Qf(\omega, t) e^{-irt} dt \\ &= \frac{1}{(2\pi)^d} [\mathcal{F}_d Qf(\omega, \cdot)](r) \end{aligned}$$

Proof. Write  $x = tw + y$  where  $t \in \mathbb{R}$  and  $y \perp \omega$ .

Then  $dx = dy dt$ . Furthermore

$$(rw) \cdot x = t$$

as  $\|\omega\| = 1$  and  $\omega \cdot y = 0$ . Thus

$$\begin{aligned} (2\pi)^{d/2} \mathcal{F}_d f(r\omega) &= \int_{-\infty}^{\infty} f(tw+y) e^{-irt} dy dt \\ &= \int_{-\infty}^{\infty} Qf(\omega, t) e^{-irt} dt. \end{aligned}$$

Corollary The Radon transform is injective on  $L^1(\mathbb{R}^d)$

Proof. It follows from the injectivity of the Fourier transform  $\blacksquare$

The central slice theorem implies also the following filtered back-projection (inversion) formula. Note, that if  $f$  is rapidly decreasing (i.e.  $f \in S(\mathbb{R}^d)$ ) then all integrals make sense. We will discuss the back-projection in more detail in the next section. For  $f \in S(\mathbb{R}^d)$  define

$$GRf(\omega, t) = \frac{1}{2} \int_{-\infty}^{\infty} \int_{\mathbb{R}^{d-1}} F_Q f(\omega, r) e^{-itr} |r|^{n-1} dr$$

1D FT in the second variable.

Then a simple calculation shows:

Theorem (Filtered Back-projection formula)

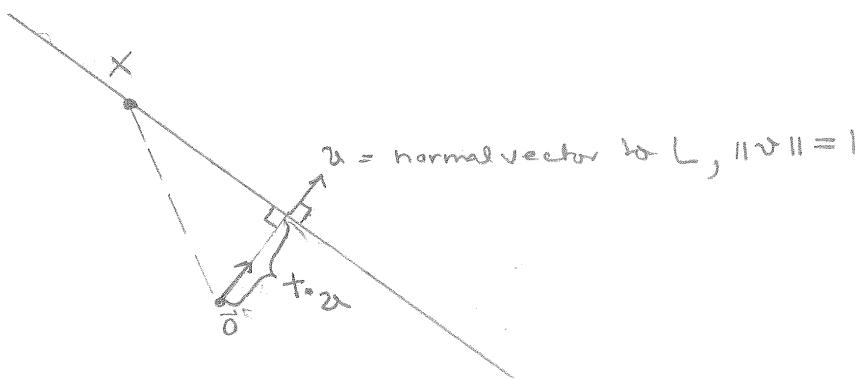
Let  $f \in S(\mathbb{R}^d)$ . Then  $\exists c$

$$f(x) = c \int_{S^{d-1}} GRf(\omega, \omega \cdot x) d\sigma(\omega).$$

#### § 4 The Back-projection

For  $x \in \mathbb{R}^d$  let  $\tilde{X}$  denote the set of hyperplanes containing  $x$ . Then

$$\tilde{X} = \{(\omega, \omega \cdot x) \mid \omega \in S^{d-1}\} \cong S^{d-1}.$$



Hence, if  $\varphi: \Xi \rightarrow \mathbb{C}$  is continuous, then we can define a function on  $\mathbb{R}^d$  by average of  $\varphi$  over all hyperplanes containing  $x$

$$R^V \tilde{\varphi}(x) = \check{\varphi}(x) = \int_{S^{d-1}} \varphi(\omega, x \cdot \omega) d\sigma(\omega)$$

The integral operator  $R^V$  is called the back-projection or dual Radon transform. We have

$$\int_{\mathbb{R}^d} f(x) \check{\varphi}(x) dx = \int_{\mathbb{R}^d} \hat{f}(\xi) \varphi(\xi) d\xi.$$

Denote by  $\Omega_k = \frac{2\pi^{k/2}}{\Gamma(k/2)}$  the area of the unit sphere in  $\mathbb{R}^k$ .

Lemma Let  $f \in S(\mathbb{R}^d)$ . Then  $\exists c_n$  such that  $\int_{\mathbb{R}^d} |x-y|^{-n} f(y) dy = c_n \int_{\mathbb{R}^d} |y|^{-n} f(y) dy$

$$(\hat{f})^V(x) = c_n \int_{\mathbb{R}^d} |x-y|^{-n} f(y) dy$$

Proof. We have

$$(\hat{f})^V(x) = \int_{S^{d-1}} \hat{f}(\omega, \omega \cdot x) d\sigma(\omega)$$

all constants are known explicitly and depend only on the normalization of the measure

$$= \iint_{S^{d-1} \times \omega^\perp} f((x \cdot \omega)\omega + y) dy d\omega$$

Next, note that  $x - (x \cdot \omega)\omega \in \omega^\perp$  as  $\omega \cdot (x - (x \cdot \omega)\omega) = \omega \cdot x - x \cdot \omega \|\omega\|^2 = 0$ . Hence

$$\begin{aligned} \iint_{S^{d-1} \times \omega^\perp} f((x \cdot \omega)\omega + y) dy d\omega &= \iint_{S^{d-1} \times \omega^\perp} f(x + y) dy d\omega \\ &= c \int_{\mathbb{R}^d} f(x+y) |y|^{-1} dy \\ &= c \int_{\mathbb{R}^d} \frac{f(y)}{|x-y|} dy. \end{aligned}$$

■

Lemma Let  $f \in \mathcal{J}(\mathbb{R}^d)$ , then

$$\mathcal{P}^\nu \mathcal{P} f(x) = c \int_{\mathbb{R}^d} f(x+y) |y|^{1-n} dy. \quad ■$$

Note the occurrence of the  $y \mapsto |y|^{-n}$  in both cases.

### § 5 Inversion formulas involving powers of the Laplacian

For  $\operatorname{Re} s < n$  define

$$I^s f(x) = \frac{1}{\pi} \int_{\mathbb{R}^d} (|\lambda|^{-s} \hat{f}(\lambda))(\lambda) d\lambda$$

$I^s$  is called the Riesz potential.

Recall, the following

$$\mathcal{F}\left(\frac{\partial f}{\partial x_j}\right)(\lambda) = i\lambda_j \mathcal{F}f(\lambda)$$

Let  $\Delta = \sum_{j=1}^d \left(\frac{\partial}{\partial x_j}\right)^2$  be the Laplacian.

Then

$$\mathcal{F}(\Delta f)(\lambda) = -|\lambda|^2 \mathcal{F}f(\lambda).$$

In particular, if  $\alpha = -2m \in -2\mathbb{N}$ , then

$$I^\alpha f(x) = (-\Delta)^m f(x)$$

and

$$I^0 f(x) = f(x).$$

We note also, that (formally)

$$I^{\alpha+\beta} = I^\alpha I^\beta.$$

Finally,  $\alpha \mapsto I^\alpha f(x)$  has a meromorphic continuation to all of  $\mathbb{C}$ .

Theorem Let  $f \in \mathcal{S}(\mathbb{R}^d)$ . Then  $\exists c \neq 0$ :

$$\begin{aligned} f &= c I^{1-d} (\hat{f})^\vee \\ &= c (-\Delta)^{\frac{d-1}{2}} (\hat{f})^\vee \end{aligned}$$

Remark: Note the difference between  $d$  even and  $d$  odd. If  $d$  is odd, then  $(d-1)/2$  is an integer and hence  $(-\Delta)^{\frac{d-1}{2}}$

is a differential operator and hence local

On the other hand, if  $d$  is even, then  $\frac{d-1}{2}$  is a half-integer and  $(-\Delta)^{\frac{d-1}{2}}$  is a pseudo-differential operator and involves global information.

There are other ways to write the inversion formula.

We can consider the 1-D Riesz potential and write

$$f = c \int_{\mathbb{R}^d} I^{1-d} \hat{f}.$$

For  $g \in S(\mathbb{R})$  we have

$$\begin{aligned} (I^{1-d} g)^\wedge(x) &= |\lambda|^{d-1} \hat{g}(\lambda) \\ &= (\text{sign}(\lambda))^{d-1} \lambda^{d-1} \hat{g}(\lambda) \\ &= (\text{sign}(\lambda))^{d-1} \widehat{((-i)^{d-1} g^{(d-1)})}(\lambda), \end{aligned}$$

where

$$\text{sign}(\lambda) = \begin{cases} 1 & \lambda > 0 \\ -1 & \lambda < 0. \end{cases}$$

Define the Hilbert transform by

$$\begin{aligned} (Hg)(x) &= \mathcal{F}^{-1}(-i \text{sign}(\lambda) \hat{g}(\lambda))(x) \\ &= c \int_{-\infty}^{\infty} \frac{h(t)}{s-t} dt. \end{aligned}$$

The inversion formula can now be rewritten as  
d even:

$$\begin{aligned}
 f(x) &= c \int_{S^{d-1}} H \hat{f}^{(d-1)}(\omega, x \cdot \omega) d\omega \\
 &= c \int_{S^{d-1}} \int_{-\infty}^{\infty} \frac{\hat{f}^{(n-1)}(\omega, t)}{x \cdot \omega - t} dt d\omega \\
 &= c \int_{S^{d-1}} \int_{-\infty}^{\infty} \frac{\hat{f}^{(n-1)}(\omega, x \cdot \omega + q)}{q} dq d\omega
 \end{aligned}$$

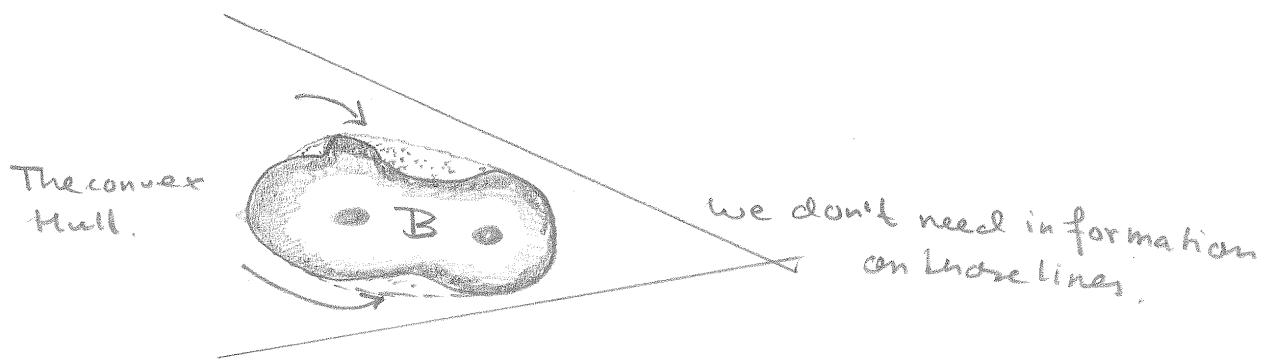
d-odd:

$$f(x) = c \int_{S^{d-1}} \hat{f}^{(n-1)}(\omega, x \cdot \omega) d\omega$$

This form of the inversion formula shows clearly the difference between the even and odd dimension. In odd dimension we only need to know  $\hat{f}$  on hyperplanes in a neighborhood of  $\tilde{x}$  whereas in even dimension we need information on  $\hat{f}$  on all hyperplanes.

## § 6 Support Theorems

X-ray tomography is the 2-D Radon transform and we have seen, that the reconstruction formula involves global information on  $\hat{f}$ . The following shows, that under the assumption  $f = 0$  outside the body of interest, we only need information on  $\hat{f}$  on lines, that actually go through the body.



Theorem Let  $f \in \mathcal{S}(\mathbb{R}^d)$  and  $C \subseteq \mathbb{R}^d$  be convex and compact. If  $\hat{f}(L) = 0$  for all lines  $L$  s.t.  $L \cap C = \emptyset$  then  $f|_{\mathbb{R}^d \setminus C} = 0$ .

Remark: It is not necessary to assume, that  $f \in \mathcal{S}(\mathbb{R}^d)$  but one has to assume that  $f$  vanishes "fast" at  $\infty$ .