

## 1 $\|P_n\|$ : The Norm of the $n$ th Legendre Polynomial

We will prove that  $\|P_n\| = \sqrt{\frac{2}{2n+1}}$  for each nonnegative integer  $n$ , leaving numerous details for the reader to check.

We will use Rodriguez's Formula<sup>1</sup>, established in class:  $P_n(x) = \frac{1}{2^n n!} \frac{d^n u}{dx^n}$ , where  $u = (x^2 - 1)^n$ . We will apply integration by parts repeatedly.

$$\begin{aligned} (2^n n!)^2 \int_{-1}^1 P_n^2 dx &= \int_{-1}^1 u^{(n)} u^{(n)} dx = u^{(n)} u^{(n-1)} \Big|_{-1}^1 - \int_{-1}^1 u^{(n+1)} u^{(n-1)} dx = \int_{-1}^1 u^{(n+1)} u^{(n-1)} dx \\ &= \dots = (-1)^n (2n)! \int_{-1}^1 (x^2 - 1)^n dx = (2n)! \int_{-1}^1 (1 - x^2)^n dx \\ &\stackrel{(A)}{=} (2n)! \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} \cos^{2n+1} \theta d\theta \stackrel{(B)}{=} (2n)! \frac{2(2^n n!)^2}{(2n+1)!} \end{aligned}$$

where step (A) is a trigonometric substitution, and step (B) is explained as follows.

We prove by induction that  $\int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} \cos^{2n+1} \theta d\theta = 2 \int_0^{\frac{\pi}{2}} \cos^{2n+1} \theta d\theta = \frac{(2^n n!)^2}{(2n+1)!}$ . The case  $n = 0$  is easy and is left as an exercise. Now assume that the formula is true for  $n - 1$  and use integration by parts to prove the validity of the formula for  $n$ .

Finally, the steps above establish the claimed norm for  $P_n$ .

## 2 Zeros of the Bessel Functions: Heuristic Explanation

Consider the Bessel equation for any  $\nu \geq 0$ :

$$x^2 y'' + xy' + (x^2 - \nu^2) y = 0.$$

Make the substitution  $y = \frac{u}{\sqrt{x}}$ , where  $u$  will be the new dependent variable, replacing  $y$ , and  $x$  remains the independent variable. Show that this transforms the general Bessel equation into

$$u'' + \left(1 + \frac{\frac{1}{4} - \nu^2}{x^2}\right) u = 0.$$

Observe that if  $x \gg 0$ , which is read as  $x$  is *very much bigger* than zero, then this equation is quite similar to  $u'' + u = 0$ . The latter equation is solved by  $\cos x$  and by  $\sin x$ , which are functions that oscillate endlessly between positive and negative values. In fact, the equation tells us that  $u$  *accelerates* toward negative values whenever  $u$  is positive, and vice versa. This suggests strongly that  $u$ , and also  $y$ , must cross the  $x$ -axis infinitely often. Details of a rigorous argument can be found in Watson's book.

## 3 Norm of the Bessel Function $J_n$ , $n \geq 0$

Note that  $n$  does not need to be an integer in this section. We will prove that

$$\|J_n(\lambda_{mn} x)\|^2 = \frac{R^2}{2} J_{n+1}^2(\lambda_{mn} R).$$

<sup>1</sup>These and other topics can be studied further in Watson's classic treatise, *Bessel Functions*.

Here  $\lambda_{mn}R = a_{mn}$ , the  $m$ th zero of  $J_n$ , and  $\|J_n(\lambda_{mn}x)\|^2 = \int_0^R J_n^2(\lambda_{mn}x) x dx$ .

Denote  $\phi_n(x) = J_n(\lambda x)$  for general  $\lambda \in \mathbb{R}$ . It follows that

$$\begin{aligned}
 [x\phi'_n(x)]' + \left(\frac{-n^2}{x} + \lambda^2 x\right) \phi_n(x) &= 0. \text{ Now multiply by } 2x\phi'_n(x) : \\
 ([x\phi'_n(x)]^2)' + (\lambda^2 x^2 - n^2) [\phi_n^2(x)]' &= 0. \text{ Next integrate both sides:} \\
 [x\phi'_n(x)] \Big|_0^R &= - \int_0^R (\lambda^2 x^2 - n^2) [\phi_n^2(x)]' dx. \quad (*)
 \end{aligned}$$

Recall the identity

$$\begin{aligned}
 [x^{-n} J_n(x)]' &= -x^{-n} J_{n+1}(x) \text{ and proceed as follows.} \\
 -nx^{-n-1} J_n(x) + x^{-n} J'_n(x) &= -x^{-n} J_{n+1}(x) \\
 -nJ_n(x) + xJ'_n(x) &= -xJ_{n+1}(x). \text{ Next replace } x \text{ by } \lambda x : \\
 \lambda x \frac{\phi'_n(x)}{\lambda} &= n\phi_n(x) - \lambda x \phi_{n+1}(x).
 \end{aligned}$$

Combining the latter calculations we find from the left side of Equation (\*) that

$$\begin{aligned}
 [x\phi'_n(x)]^2 \Big|_0^R &= [n\phi_n(x) - \lambda x \phi_{n+1}(x)]^2 \Big|_0^R \text{ (Now set } \lambda = \lambda_{mn}) \\
 &= \lambda_{mn}^2 R^2 J_{n+1}^2(\lambda_{mn}R) \text{ since } J_n(\lambda_{mn}R) = 0 \text{ if } n \geq 1. \text{ Integrate (*) on the right by parts:} \\
 &= - \underbrace{(\lambda_{mn}^2 x^2 - n^2) J_n^2(\lambda_{mn}x)}_{=0} \Big|_0^R + 2\lambda_{mn}^2 \int_0^R x J_n^2(\lambda_{mn}x) dx.
 \end{aligned}$$

It follows that  $\int_0^R J_n^2(\lambda_{mn}x) x dx = \frac{R^2}{2} J_{n+1}^2(\lambda_{mn}R)$ , which concludes the proof.